DARK MATTER THEORY

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IV International Pontecorvo Neutrino Physics School Alushta, Crimea (Ukraine) – 30.09.2010

Outline of the lectures

- Evídence of dark matter
- Where does it come from?
- What is made of? Any clue on or relation with New Physics?
- How is it distributed aroud us?
- Can we directly or indirectly detect its presence in our (close) astrophysical environment?



Geometry: the Universe is Flat Dynamics: the Universe is expanding Decelerate for most of its history
 Accelerate since "recent" time and at very "old" times (inflation)

- Ω_{T} CMB temperature anisotropies
- Ω_{Λ} Luminosity distance of high-z SNIa
- Ω_{M} Clustered mass abundance
- Primordial Nucleosynthesis Amplitude of CMB temperature anisotropies $\Omega_{
 m B}$



Dark Matter

Dynamics of galaxy clusters Rotational curves of galaxies Weak lensing Structure formation from primordial density fluctuations Energy density budget









Dark Matter







Dynamics of galaxy clusters Rotational curves of galaxies Weak lensing Structure formation from primordial density fluctuations Energy density budget

Solution to the DM problem

Gravitational: gravity is not of Einstein form on large scales

Particle Physics: DM is of particle nature



Dynamics of galaxy clusters Rotational curves of galaxies Weak lensing Structure formation from primordial density fluctuations Energy density budget

(*) Standard neutrino: Too light: act as HDM (not CDM)



Particle candidate

Cosmology of the particle DM

Relic from the early U.

Astrophysical signals of the particle DM

Detection

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COSMOLOGY OF THE PARTICLE DM



Particle DM thermalization in the early Universe

Equilibrium reactions

$$\chi A \longleftrightarrow \chi A$$
$$\chi \bar{\chi} \longleftrightarrow \bar{A}A$$

	Relativistic Bosons	Relativistic Fermions	Non-relativistic (Either)
n_i	$\frac{\zeta(3)}{\pi^2}g_iT^3$	$\left(\frac{3}{4}\right)\frac{\zeta(3)}{\pi^2}g_iT^3$	$g_i \left(\frac{m_i T}{2\pi}\right)^{3/2} e^{-m_i/T}$
ρ_i	$\frac{\pi^2}{30}g_iT^4$	$\left(\frac{7}{8}\right)\frac{\pi^2}{30}g_iT^4$	$m_i n_i$
p_i	$\frac{1}{3}\rho_i$	$\frac{1}{3} ho_i$	$n_i T \ll \rho_i$

$$\Gamma = n \langle \sigma v \rangle$$
: interaction rate

$$\langle \sigma v \rangle = \frac{\int d^3 p_i \ d^3 p_j \ f_i(E) \ f_j(E) \ \sigma_{ij} v_{ij}}{\int d^3 p_i \ d^3 p_j \ f_i(E) \ f_j(E)}$$

 $H = \dot{a}/a$: expansion rate





Neutrínos as HDM: relic abundance

$$\Omega_{\nu}h^2 = \frac{\sum_i m_i}{93 \text{ eV}}$$

$$\Omega_{\nu}h^2 \le (\Omega_{\rm DM}h^2) = 0.13 \qquad \longrightarrow \qquad \sum_i m_i \le 12 \text{ eV}$$

Neutrínos as HDM





CDM relic abundance

• Boltzmann equation:

$$\frac{dn}{dt} = -3Hn - \langle \sigma v \rangle (n^2 - n_{\rm eq}^2)$$

$$\frac{\mathrm{d}Y}{\mathrm{d}x} = \beta \frac{g_{*S}}{g_{*}^{1/2}} < \sigma_{\mathrm{ann}} v_r > [Y^2 - Y_{eq}^2]$$
$$Y = \frac{n}{s} \qquad x = T/m_{\chi} \qquad \beta = 0.264m_{\chi} M_P$$

• <u>Relic abundance</u>:

$$\Omega_{\chi} h^2 = 8.5 \cdot 10^{-11} \frac{g_*^{1/2}(x_f)}{g_{*S}(x_f)} \left(\frac{\text{GeV}^{-2}}{<\sigma_{\text{ann}} v_r >_{\text{int}}} \right)$$

$$x_f^{-1} = \ln \left[0.145 \beta \frac{g}{g_{\star}^{1/2}} \ x_f^{1/2} \langle \sigma_{\rm ann} v_r \rangle_{(x_f)} \right] \qquad \text{freeze-out temperature}$$

The "WIMP" miracle

WIMP: Weakly Interacting Massive Particle

$$m_{\chi} \sim (\text{GeV} \div \text{TeV})$$

 $\langle \sigma v \rangle = G_F^2 m_{\chi}^2$
naturally $\Omega_{\chi} h^2 \sim 0.1$
 $x_F \equiv m_{\chi}/T_F \sim 15 \div 25$

Succesfull DM candidate

- Needs to be produced in the early Universe
 - Thermal relic
 - Non-thermal relic
- Needs to be "cold" (or, at least, "warm" enough)
 - For thermal production: weakly interacting and massive (WIMP)

$$\Omega h^2 \sim \langle \sigma v \rangle_{\rm ann}^{-1} \longrightarrow \langle \sigma v \rangle_{\rm ann} = 3 \cdot 10^{-26} {\rm cm}^3 {\rm s}^{-1}$$

unless coannihilation occurs

- If light, it nevertheless needs to act as "cold"

- Needs to be neutral
- Needs to be stable (absolutely, or on cosmologícal tíme scales)

DM AND NEW PHYSICS MODELS

• <u>"Mínimal" candidates</u>

- Neutríno: standard, ríght-handed, (...)
- Additional gauge multiplets (MDM), (...)
- Axíon

• <u>"Minimal" candidates</u>

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•	Su	persu	mmetríc	candí	dates

- Neutralino
- Sneutríno
- Gravítíno
- Axíno

SUP.	ERSYMMET	RY: I	ERMION	<u> </u>	\rightarrow B080	N
Normal particles	/fields	Supersyn	nmetric partne	ers		
		Interactio	on eigenstates		Mass eige	enstates
Symbol	Name	Symbol	Name		Symbol	Name
q = d, c, b, u, s, t	quark	\tilde{q}_L, \tilde{q}_R	squark		\tilde{q}_1,\tilde{q}_2	squark
$l = e, \mu, \tau$	lepton	\tilde{l}_L, \tilde{l}_R	slepton		\tilde{l}_1, \tilde{l}_2	slepton
$\nu = \nu_e, \nu_\mu, \nu_\tau$	neutrino	$\tilde{ u}$	sneutrino		$\tilde{\nu}$	sneutrino
g	gluon	\tilde{g}	gluino		\tilde{g}	gluino
W^{\pm}	W-boson	\tilde{W}^{\pm}	wino)		
H^-	Higgs boson	\tilde{H}_1^-	higgsino	Ş	$\tilde{\chi}_{1.2}^{\pm}$	chargino
H^+	Higgs boson	\tilde{H}_2^+	higgsino	J	-,-	
В	B-field	<i></i> Ĩ	bino)		
W^3	W^3 -field	\tilde{W}^3	wino			
H_{1}^{0}	Higgs boson	$\tilde{rr0}$		->	$\tilde{\chi}^{0}_{1,2,3,4}$	neutralino
$H_2^{\hat{0}}$	Higgs boson	\tilde{H}_1^0	higgsino			
$H_3^{ ilde{0}}$	Higgs boson	H_{2}^{0}	higgsino)		

SUSY breaking \longrightarrow massive SUSY partners



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• <u>"Minimal" candidates</u>

- Neutríno: standard, ríght-handed, (...)
- Additional gauge multiplets (MDM), (...)
- Axion



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- <u>"Minimal" candidates</u>
 - Neutríno: standard, ríght-handed, (...)
- WIMP Additional gauge multiplets (MDM), (...)
 - Axion
- <u>Supersymmetric candidates</u>
- WIMP Neutralino
- WIMP Sneutríno
 - Gravítíno
- WIMP Axíno (...)
 - Extra-dimensions
- WIMP Kaluza-Klein fields (...)

• <u>"Minimal" candidates</u>

- Neutríno: standard, RH MeV, (...) Non WIMP
- Additional gauge multiplets (MDM), (...)
- Axion

Non WIMP

- <u>Supersymmetric candidates</u>
 - Neutralino
 - Sneutríno
 - Gravitino
 - Axíno (...)
 - Extra-dimensions
 - Kaluza-Klein fields

• <u>"Minimal" candidates</u>

- Neutríno: standard, RH MeV, (...)
- Additional gauge multiplets (MDM), (...)
- Axíon (...)

(...)

Supersymmetric candidates
Neutralino
Sneutrino
Gravitino
Axino
 (...)
Extra-dimensions
Kaluza-Klein fields



• <u>"Minimal" candidates</u>

- Neutríno: standard, RH MeV, (...)
- Additional gauge multiplets (MDM), (...)
- Axion
- Supersymmetric candidates
 Neutralino
 Sneutrino
 Gravitino
 Axino
 (...)
 Extra-dimensions
 Kaluza-Klein fields
 (...)



Neutralino annihilation cross section

 $\langle \sigma_{
m ann} v
angle$



*	$\chi\chi \to f\bar{f}$	$\begin{array}{c} Z, h, H, A \\ \tilde{f}_L, \tilde{f}_R \end{array}$	s channel t and u channels
*	$\chi\chi \to hh, hH, HH, AA$	$\begin{array}{c}h,H\\\chi_i\ (i=1,2,3,4)\end{array}$	s channels t and u channels
*	$\chi\chi \to hA, HA$	Z, A $\chi_i \ (i = 1, 2, 3, 4)$	s channel t and u channels
*	$\chi\chi \to H^+ H^-$	Z, h, H $\chi_j^+ \ (j = 1, 2)$	s channel t and u channels
*	$\chi\chi \to ZZ$	$\begin{array}{c}h,H\\\chi_i\ (i=1,2,3,4)\end{array}$	s channel t and u channels
*	$\chi\chi \to W^+W^-$	Z, h, H $\chi_j^+ \ (j = 1, 2)$	s channel t and u channels
*	$\chi\chi \to hZ, HZ$	Z, A $\chi_i \ (i = 1, 2, 3, 4)$	s channel t and u channel
*	$\chi\chi \to AZ$	$\begin{array}{c}h,H\\\chi_i\ (i=1,2,3,4)\end{array}$	s channel t and u channels
*	$\chi\chi \to W^\pm H^\mp$	$\begin{array}{c} h, H, A \\ \chi_j^+ \ (j=1,2) \end{array}$	s channel t and u channels

which one is open depends on the neutralino mass

MSSM



Minimal Supergravity



 $M_{1/2} m_0 A_0 \tan\beta \operatorname{sign}(\mu)$

H. Baer, 0901.4732 [hep-ph]

Galactic Dark Matter

CDM in galaxies:

- DM as a non-baryonic particle
- Massive particle with weak-type interactions (WIMP)
- Distributed to form a halo
 - Thermal component
 - Substructures
 - Non-thermal component

Galactic dark matter detection:

- Identify types of signals
- Exploit specific signatures
- Exploit (anti) correlations among signals
- Study relevant backgrounds
- Quantify uncertainties



DM DISTRIBUTION IN GALAXIES
Galactic environment





From numerical simulations

Navarro et al., arXív:0810.1522

Vogelsberger et al., arXív:0812.0362

Subhalos



The Aquarius Project



Most subhalos are in the outer halo

Springel et al., MNRAS 391 (2008) 1685

Velocity streams



Vogelsberger et al., arXív:0812.0362

"Canonical" halo

$$\rho(r) \longrightarrow \rho_0 = 0.3 \text{ GeV cm}^{-3}$$

$$\rho(r) \longrightarrow r^{-1} [r \to 0]$$

Recent determinations [1-3]

$$\rho_0 = 0.385 \pm 0.027 \text{ GeV cm}^{-3} \quad \text{(Einasto)}$$

$$\rho_0 = 0.389 \pm 0.025 \text{ GeV cm}^{-3} \quad \text{(NFW)}$$

[2] $\rho_0 = 0.43(11)(10) \text{ GeV cm}^{-3}$

$$f(\vec{v}) = N \exp(-v^2/v_0^2)|_{v_{\rm esc}}$$
$$v_0 = (220 \pm 50) \text{km s}^{-1}$$
$$v_{\rm esc} = (450 \div 650) \text{ km s}^{-1}$$

Streams may have impact Anisotropies may be present

[1] Catena, Ullio, arXiv:0907.0018
 [2] Salucci et al. arXiv:1003.3101
 [3] Pato et al., arXiv:1006.1322

ASTROPHYSICAL SIGNALS

MultiChannel search of WIMP dark matter

• <u>Direct search</u>: elastic scattering of χ off nuclei in a low background detector

 recoil energy of the nucleus annual modulation of the rate directionality of the recoil

- signals due to $\chi\chi$ annihilation taking place inside celestial bodies (Sun, Earth) where χ have been captured and accumulated

Neutrino flux — up-going muons in a neutrino telescope source location/some spectral feature

- signals due to $\chi\chi$ annihilation taking place in the galactic halo



Indírect searches:



Relic abundance Indirect signals



Direct detection Neutrinos from Earth and Sun



Accelerator searches

DIRECT DETECTION

Direct detection



Interaction mechanisms - WIMPs

• Elastic scattering with nuclei $\chi \: \mathcal{N} \longrightarrow \chi \: \mathcal{N}$ - Ex.: Neutralino, Sneutrinos, KK

$$E_R = \mu_N^2 v^2 (1 - \cos \theta) / m_N$$

$$E_R > few KeV$$

• Inelastic scattering with nuclei Tucker-Smith, Weiner, PRD 64 (2001) 043502

 $\chi \mathcal{N} \longrightarrow \chi' \mathcal{N}$

- Ex.: Sneutrínos

Scatter if:
$$\Delta m < \frac{\beta^2 m_1 m_N}{2(m_1 + m_N)}$$

about 1-100 KeV

Interaction mechanisms - non WIMPs

• Inelastíc, scatter on electrons

- Ex.: Light (KeV) [pseudo]scalars





WIMPs - Scattering cross section

• Spin-independent

Cross section proportional to the (mass number)² of the nucleus
 Channels:

- > Vector boson (Z)-mediated: gauge-type, well known
- > Scalar (H, squarks)-mediated: large hadronic uncertainties

$$\begin{split} I_{h,H} &= \sum_{q} k_{q}^{h,H} m_{q} \langle N | \bar{q}q | N \rangle = k_{u-\text{type}}^{h,H} g_{u} + k_{d-\text{type}}^{h,H} g_{d} \\ \hline \frac{(\text{in MeV})}{\text{Set A}} &= \frac{N | \bar{q}_{l}q_{l} | N \rangle}{27} & \frac{M | \bar{s}s | N \rangle}{131} & \frac{M | \bar{h}h | N \rangle}{56} & \frac{g_{u}}{139} & \frac{g_{d}}{214} \\ \hline \text{Set B} &= 28 & 186 & 52 & 132 & 266 \\ \hline \text{Set C} &= 37 & 456 & 30 & 97 & 523 \end{split}$$

- Nuclear form factors $F(E_R)$

• Spín-dependent

- Cross section proportional to the $(spin)^2$ of the nucleus - Spin form factors $S(E_R)$

Example: Neutralino-quark scattering





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Interaction rate (WIMP; scalar interaction)

$$\frac{dR}{dE_R} = N_T \frac{\rho_0}{m_\chi} \frac{m_N}{2\mu_1^2} A^2 \left[\xi \sigma_{\text{scalar}}^{(\text{nucleon})}\right] F^2(E_R) \mathcal{I}(v_{\text{min}})$$

$$\mathcal{I}(v_{\min}) = \int_{w \ge v_{\min}} d^3 w \;\; rac{f_{\mathrm{ES}}(ec{w})}{w}$$

$$f_{\rm ES}(\vec{w}) = f(\vec{w} + \vec{v}_{\oplus})|_{[v_{\rm rot}; v_{\rm esc}]}$$
$$v_{\rm min} = [m_N E_R / (2\mu_A^2)]^{1/2}$$

Local motions



Interaction rate (WIMP; scalar interaction)

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Response function



A. Bottíno, F. Donato, N. Fornengo, S. Scopel, PRD 72 (2005) 083521

Differential Rate - Energy Dependence



Local motions



Typical signatures of direct detection



Annual Modulation of the rate

$$\frac{dR}{dE_R}[\eta(t)] = \frac{dR}{dE_R}[\eta_0] + \frac{\partial}{\partial\eta} \left(\frac{dR}{dE_R}\right)_{\eta=\eta_0} \Delta\eta \,\cos[\omega(t-t_0)]$$

$$= S_0(E_R) + S_m(E_R) \cos[\omega(t-t_0)]$$

$$\eta(t) = v(t)/v_0$$

Annual modulation



N. Fornengo, S. Scopel, PLB 576 (2003) 189



Current direct detection experiments

- Background-rejection experiments (CDMS, XENON, CoGeNT, +)
 - Do not exploit a specific signature of the signal
 - Rely on reduction/interpretation of background

- Annual modulation experiments (DAMA)
 - Exploit a specific signature
 - Required to be highly stable over long periods

DAMA annual modulation

Effect at 8.9 σ C.L.

Single-hit events in the signal energy-window Stability parameters do not modulate

Compatible to DM scatter offnuclei on electrons

Cumulative exposure: 1.17 ton x yr (13 annual cycles) (i.e. 427050 Kg x day)

 $S_m[2-6 \text{ KeV}] \approx (0.0116 \pm 0.0013) \text{ cpd/kg/keV}$

Phase = (146 ± 7) days

Period = (0.999 ± 0.002) years





2-6 keV





DAMA annual modulation region



A. Bottino, F. Donato, N. Fornengo, S. Scopel, PRD 78 (2008) 083520

CDMS II – Fínal exposure (2009) - 2 events pass cuts



Cumulative exposure: 612 kg x day

Z. Ahmed (CDMS Collab.), arXiv:0912.3592 [astro-ph.CO]

CDMS 2009



A. Bottíno, F. Donato, N. Fornengo, S. Scopel, PRD 81 (2010) 107302

CoGeNT



XENON 100



XENONIOO Collaboration, arXiv:1005.0380v2 [astro-ph.CO] XENON100 Collaboratio, arXiv:1005.2615v1 [astro-ph.CO] E. Aprile, WONDER Workshop, LNGS, March 2010

Xenon detector: scintillation efficiency



XENON100 Collaboration, arXiv:1005.2615v1 [astro-ph.CO]

XENON 100



XENON100 Collaboration, arXiv:1005.0380v2 [astro-ph.CO] XENON100 Collaboratio, arXiv:1005.2615v1 [astro-ph.CO] E. Aprile, WONDER Workshop, LNGS, March 2010
CRESST - Unpublished



Seidel, CRESST Collaboration, WONDER Workshop, LNGS , March 2010 Jochum, GGI, May 2010

SEARCHES AT NEUTRINO TELESCOPES

Neutrinos from Earth and Sun

• Capture:

- galactic DM particles that cross the Earth and the Sun, can interact with the nuclei in these bodies and loose enough energy to remain gravitationally captured

• Accumulation:

- after subsequent interactions they tend to drop into the innermost parts of the Earth and the Sun, where they accumulate

• Annihilation:

- when the energy density in the inner parts of the Earth and the Sun increases enough, they may start to annihilate

neutrino flux



Capture Rate

- Elastic scattering of the DM particle with a nucleus *i* in a spherical shell at a distance *r* from the center of the Earth (or Sun)
- In order to be captured, the velocity of the DM particle after the interaction must be smaller than the escape velocity at the shell

$$v_{\rm esc}^{\rm Sun} = 618 \ {\rm Km \ s^{-1}}$$
 at the surface
 $v_{\rm esc}^{\rm Earth} = 11.2 \ {\rm Km \ s^{-1}}$ mean DM particle velocity

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Capture Rate

$$C = \sum_{i} \left(\frac{8}{3\pi}\right)^{1/2} \left[\sigma_{i} \frac{\rho_{\chi}}{m_{\chi}} \bar{v}\right] \left[\frac{M_{i}}{m_{i}}\right] \left[\frac{3v_{esc}^{2}}{2\bar{v}^{2}} \langle \phi \rangle_{i}\right] \xi(\infty) S_{i}$$

Sun

Nuc	Н	He	0	С	Ne	Fe	N	Sí	Mg
f	0.77	0.21	8.10-3	4.10 ⁻³	1.10-3	1.10-3	9.10-4	8.10-4	7.10-4
А	1	4	16	12	20	56	14	28	24

nuclei of type *i* in the body





Capture rate on Earth Sun



J. Lundberg, J. Edsjo, astro-ph/040111

Accumulation and concentration

- DM particles which have been captured inside Earth or Sun can suffer subsequent scatterings
- This may lead to:
 - Concentration in the innermost parts of the Earth or Sun
 - Development of an equilibrium distribution of these particles

distribution
$$n(r) = n_0 e^{-\alpha_B m_{\chi} r^2}$$

 n_0 central density $\alpha_B = 2\pi G \rho_0 / (3T_0)$

Annihilation rate

$$\Gamma_A = \frac{C}{2} \tanh^2\left(\frac{t_0}{\tau_A}\right)$$

Capture rate CAge of the body $\tau_0 = 4.6 \text{ Gyr}$ Relaxation time $\tau_A = [CC_A]^{-1/2}$ $C_A = \langle \sigma_{ann} v \rangle_0 V_2 / V_1^2$ $V_j = c_B (jm_\chi/10 \text{ GeV})^{-3/2} \text{ cm}^3$ Effective volumes of DM concentrations More concentrated for larger masses

$$c_B = 1.8 \cdot 10^{25} \ / \ 6.6 \cdot 10^{28}$$

Earth Sun

Neutrino Production

- Neutrinos are produced by DM annihilation
 - Available channels depend on mass threshold $\chi \chi \rightarrow \nu \nu, l\bar{l}, q\bar{q}, W^+W^-, ZZ$, Higgses, Higgs + gauge

- Quark hadronize \rightarrow neutrinos from hadron decay

- Productions in Earth
 - Muons: stopped before decay \rightarrow neutrinos below typical thresholds
 - Taus: decay almost as in vacuum
 - Light hadrons: typically stopped before decay
 - Heavy hadrons: typically decay before loosing significant energy
- Production in Sun
 - Leptons: stopping power of medium is stronger \rightarrow softer neutrino spectra
 - Light hadrons: typically stopped before decay
 - Heavy hadrons: energy losses important, need modeling

Spectra at production



M. Cirelli, N.Fornengo, T. Montaruli, I. Sokalski, A. Strumia, F. Vissani, NPB 727 (2005) 99

Neutrino Propagation

Density matrix evolution



Effect of propagation



Earth:

– Affected only by "atmospheric" oscillation $\nu_{\mu} \leftrightarrow \nu_{\tau}$ at E < 100 GeV

Sun:

- Affected by average "solar" and "atmospheric" oscillations
- Absorption suppresses neutrinos for E > 100 GeV (partially converted to lower energy neutrinos (by NC and regeneration)

M. Círellí, N.Fornengo, T. Montarulí, I. Sokalskí, A. Strumía, F. Vissaní, NPB 727 (2005) 99 See also: M. Blennow, J. Edsjo, T. Ohlsson, JCAP 0801 (2008) 021 for an event-based MC approach

Earth signal: through-going muons



Neutralino DM in Gaugino non-universal MSSM

 m_{χ} (GeV)

V. Níro, A. Bottíno, N.Fornengo, S. Scopel, arXív:0909.2348

 m_{χ} (GeV)

Sun signal: stopping muons



V. Níro, A. Bottíno, N.Fornengo, S. Scopel, arXív:0909.2348

Sun signal: muons above I GeV



G. Wilkstrom, J. Edsjo, arXiv:0903.2986

Reconstruction of the DM properties



M. Círellí, N.Fornengo, T. Montarulí, I. Sokalskí, A. Strumía, F. Vissaní, NPB 727 (2005) 99

Any effect of trapped DM on Sun's properties?

• WIMPs are confined in the central region of the Sun

 $R_{\chi} \simeq 0.01 \, (100 \, {
m GeV}/m_{\chi})^{1/2} \, R_{\odot}$

• They provide a mechanism for energy transport



Maximal effects



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Neutralino DM in the MSSM



ANTIPROTONS ANTIPROTONS





- <u>Physical processes</u>
 - Diffusion: uniform in the whole (disk + diffusive halo) volume
 - Inelastic (non-annihilating) scattering and annihilation
 - Galactic wind away from the disk in vertical direction
 - Energy losses:
 - \checkmark Ionization: interaction with the neutral IS matter
 - ✓ Coulomb scattering: interaction with ionized plasma (thermal electrons)
 - Reacceleration on random hydrodynamic waves (in the disk only)

 $\begin{array}{c} \begin{array}{c} Propagation in the Galaxy & \Phi^{\bar{p}}(r,z,T_{\bar{p}}) \end{array} \\ \hline \Phi^{\bar{p}}(r,z,T_{\bar{p}}) & \bullet^{\bar{p}}(r,z,T_{\bar{p}}) \end{array} \\ \hline \bullet \ solution of the steady-state diffusion equation with energy losses and reacceleration & \bullet \ energy losses and reacceleration$

The params are constrained by stable nuclei propagation, mainly B/C [D. Maurin et al. Astron. Astrophys. 381 (2002) 539]

case	δ	K_0	L	V_c	V_A	$\chi^2_{ m B/C}$
		$(\rm kpc^2/Myr)$	(kpc)	$(\rm km/sec)$	$(\rm km/sec)$,
max	0.46	0.0765	15	5	117.6	39.98
med	0.70	0.0112	4	12	52.9	25.68
min	0.85	0.0016	1	13.5	22.4	39.02



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Secondary antiprotons



Antiproton flux

Antiproton/proton fraction

F. Donato, D. Maurín, P. Brun, T. Delahaye, P. Salatí, arXív.0810.5292 [astro-ph]



A. Bottíno, F. Donato, N.F., S. Scopel, PRD 70 (2004) 015005

F. Donato, N.F., D. Maurín, P. Salatí, R. Taíllet, PRD 69 (2003) 063501

Sneutrinos in Left-Right models



C. Arína, N. Fornengo, JHEP 0711 (2007) 029

ANTIDEUTERONS ANTIDEUTERONS ANTIDEUTERONS





- Tertiary redistribution

A. Donato, N. Fornengo, D. Maurín, PRD 78 (2008) 043506

TOA fluxes and S/B gain



Signal with uncertainty band for:

- 50 GeV WIMP mass
- WMAP relic abundance

Signal/(Back+Signal) ratio

A. Donato, N. Fornengo, D. Maurín, PRD 78 (2008) 043506

Theoretical predictions MSSM + gaugino non universal SUGRA 10^{-25} 10-25 GAPS ULDB 100 events 10 even 1 event 10-26 10-26 $\xi^2 < \sigma v > (cm^3 s^{-1})$ $\xi^2 < \sigma v > (cm^3 s^{-1})$ 10-27 10-27 10 events 10-28 10-28 GAPS ULDB 1 event 10-29 10-29 100 10 100 m_{χ} (GeV) m_{γ} (GeV)

cosmologically dominant neutralinos
cosmologically subdominant neutralinos

 $0.095 \le \Omega_{\chi} h^2 \le 0.131$ $\Omega_{\chi} h^2 < 0.095$

A. Donato, N. Fornengo, D. Maurín, PRD 78 (2008) 043506

ANTIMATTER IN COSMIC RAYS POSITRONS LOCILLOUS
PAMELA positron fraction



O. Adriani et al. (PAMELA Collab.), Nature 458 (2009) 607

FERMI electrons + positrons



A. Abdo et al. (FERMI/LATCollab.)PRL 102 (2009) 181101

Leptonic "anomaly"

- "Excess" above few tens of GeV
- Hard to reconcile with pure secondary origin
- Leptonic primary sources
 - Pulsars
 - > Purely leptonic production (no protons/antiprotons)
 - Supernova remnants
 - CR sources

- Dark matter annihilation (or decay)

Primary Electrons and Positrons



J. Lavalle, T. Delahaye, R. Líneros, F. Donato, N. Fornengo, arXív:1002.1910 [astro-ph.HE]



IV Pontecorvo School - Alushta - 30.09.2010

• For astrophysical interpretation, see also:

- Hooper, Blasi, Serpico, arXiv:0810.1527
- Yuksel, Kístler, Stanev, arXív:0810.2784
- Profumo. arXív:0812.4457
- Grasso et al. arXív:0905.0636
- Malyshev, Cholís, Gelfand, arXív:0903.1310
- Kawanaka, loka, Nojiri, arXiv:0903.3782
- (...; incomplete list)



T. Delahaye, R. Líneros, F. Donato, N. Fornengo, P. Salatí, Phys. Rev. D 77 (2008) 063527







Fit on positron and antiproton data (with S&M background)

M. Cirelli, M. Kadastik, M. Raidal, A. Strumia, arXiv:0809.2409v3 [hep-ph] See also: V. Barger, W.-Y. Keung, D. Marfatía, G. Shaughnessy, arXiv:0809.0162v2 [hep-ph]

Model independent analysis





Astrophysical boost



J. Lavalle, Q. Yuan, D. Maurín, X.J. Bí, A&A 479 (2008) 427

Particle physics boost: Sommerfeld effect



M. Lattanzí, J. Sílk, arXív:0812.0360v1 [astro-ph]

See also:

J. Hisano, M. Nagai, M. Nojiri, M. Senami, PRL 92 (2004) 031303 J. Hisano, S. Matsumoto, M. Nojiri, S. Saito, PRD, 71 (2005) 063528 M. Cirelli, A. Strumia, M. Tamburini, NPB 787 (2007)

J. March-Russell, S. M. West, D. Cumberbatch, D. Hooper, JHEP 0807 (2008) 058

N. Arkani-Hamed, D. P. Finkbeiner, T. Slatyer, N. Weiner, arXiv:0810.0713 [hep-ph]

M. Čírellí, M. Kadastík, M. Raídal, A. Strumía, arXív:0809.2409v3 [hep-ph]

Astrophysical bounds on Sommerfeld boost



R. Catena, N.Fornengo, M. Pato, L. Pieri, A. Masiero, arXiv:0912.4421 astro-ph.CO]

• For bounds on enhanced annihilation rate, see also:

- S. Gallí, F. Iocco, G. Bertone, A. Melchiorri, arXiv:0905.0003v1 [astro-ph]
- M. Pato, L. Píerí, G. Bertone, 0905.0372vl [astro-ph.HE]
- Cirellí, Pancí, Serpico, arXiv:0912.0663
- Círellí, locco, Pancí, arXiv:0907.0719
- Cirelli, Panci, arXiv:0904.3080
- Bertone, Círellí, Strumía, Taoso, arXiv:0811.3744
- Círellí, Strumía, arXív:0808.3867
- Círellí, Franceschíní, Strumía, arXív:0802.3378
- Arína et al. arXív:1004.0645
- Gallí et al., arXív:1005.3808
- Feng, Manoj, Hai-Bo, arXiv:1005.4678 (...; incomplete list)

• For DM interpretation of the PAMELA data:

- More than 150 papers

"Cosmological boost"



Cosmologícal boost for PAMELA



$$H(T) = A(T)H_{\rm GR}(T)$$

 $A(T) = 1 + \eta \left(\frac{T}{T_{\rm f}}\right)^{\nu} \tanh\left(\frac{T - T_{\rm re}}{T_{\rm re}}\right)$

R. Catena, N.Fornengo, M. Pato, L. Píerí, A. Masíero, arXív:0912.4421 astro-ph.CO]

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IV Pontecorvo School - Alushta - 30.09.2010

Interpretation of leptonic CR data

• DM: problematic

- Requires large boosts
 - Astrophysical: quite unlikely
 - Particle physics (Sommerfeld): somehow contrived, constrained
 - Cosmological: constrained, requires modified cosmology
- Requires leptophilic DM: may be arranged, but not viable for most of the "canonical" DM candidate (neutralinos, sneutrinos)

• Astrophysical interpetation

- Pulsars and SNR may account for the excess
- Energetics not fully understood, but consistent with models





Subhalos and extragalactic



L. Pieri et al., arXiv:0908.0195 [astro-ph.HE]

FERMI LAT data on gamma rays



Abdo et al. , PRL 103 (2009) 251101

Bounds on cosmological DM annihilation



Abdo et al. , JCAP04 (2010) 014

Sneutrino Dark Matter



C. Arína, N. Fornengo, JHEP 0711:029,2007

Gamma-ray line

$$\begin{array}{c} \chi\chi \longrightarrow \gamma\gamma \\ \chi\chi \longrightarrow \gamma Z \end{array}$$

Typically suppressed (1-loop process)



Gamma-ray líne: neutralino DM



Gamma-ray line from neutralino DM may be visible, under favorable conditions

G. Zaharijas, D. Hooper, PRD 73 (2006) 103501

Further topics

- Neutrinos as DM messengers from the Galactic Center
 - Difficult to detect (correlate with gamma-rays)
- Radio emission from electron emission in magnetic fields
 - WMAP haze: excess of microwave emission at GC [?]
 - Spherical, radius 4 Kpc
 - Synchrotron emission from electron component?
- Gamma rays from IC of electrons/positrons on radiation fields
 - FERMI haze: excess of gamma-ray emission at GC [?]
 - Inverse Compton counterpart of the WMAP haze?
- Sunyaev-Zeldovich effect on CMB in galaxy clusters
- Colafrancesco, AA 422 (2004) L23

Finkbeiner et alrarXiv:0910.4583

- Very small effect, prospects for the future (?)

Finkbeiner, Ap. J. 614 (2004) 186

Conclusions

- Astropysical searches may be proficiently used to set constraints on the properties of particle DM
- If a signal is detected, it may guide us toward the properties of the DM candidate (and to some extent of the underlying New Physics model)
- Different detection signals probe different properties of the DM particle and feel different features of the galactic environment
- DM searches require:
 - To exploit specific and typical signatures of the various types of signals
 - Better knowledge of the astrophysical environment

Conclusions

- Cosmological properties and astrophysical signals of particle DM candidates can either guide or complement accelerator physics searches
- Viceversa, accelerators, with their capability of identifying (at least part of the) BSM particles and their properties, will allow to shape out the predictions for DM signals
- The two approaches are therefore both fundamental in the study of the DM hypothesis
 - Accelerators: prove the existence of Physics BSM and directly discover the new physical states
 - DM searches: prove that the new physical states explain the DM puzzle and explicitely identify the DM presence in the astrophysical environment
- The interplay between the two approaches may be able (in the future) to tell us something on the cosmological evolution of the early Universe

The Particle Dark Matter Crossroad

