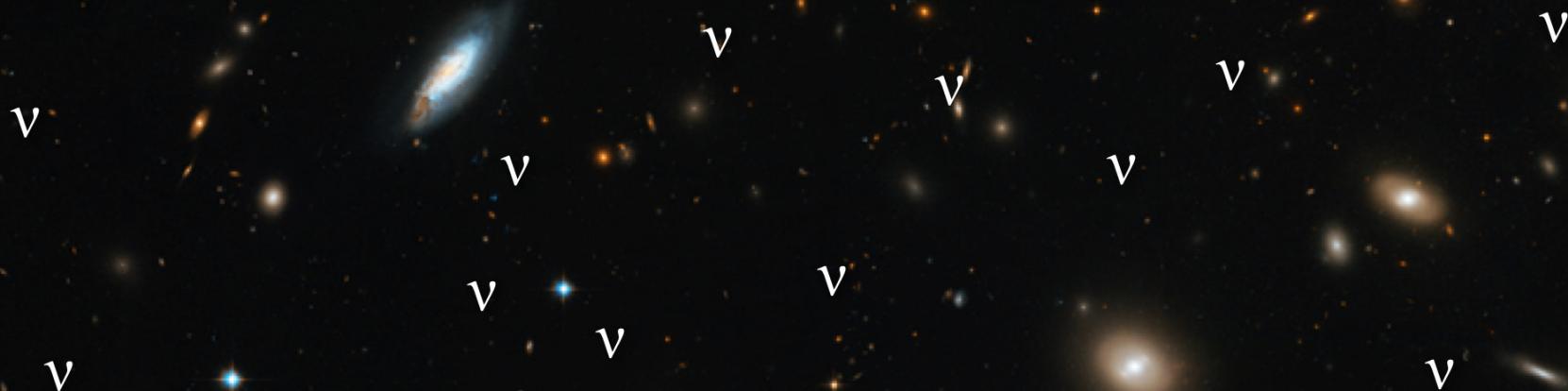


# Light neutrinos in Cosmology (II)



Picture from Hubble ST



Sergio Pastor  
(IFIC Valencia)

IV Int. Pontecorvo Neutrino  
Physics School  
Alushta, October 2010



# Light neutrinos in Cosmology

## 2nd lecture

Massive neutrinos as Dark Matter

Effects of neutrino masses on cosmological observables

Bounds on  $m_\nu$  from CMB, LSS and other data

Future sensitivities on  $m_\nu$  and  $N_{\text{eff}}$  from cosmology

# Neutrinos as Dark Matter

# History of the Universe

Accelerators: CERN-LHC  
FNAL-Tevatron  
BNL-RHIC  
CERN-LEP  
SLAC-SLC  
high-energy cosmic rays

BIG BANG

Inflation

Relativistic neutrinos

Key:	
q quark	W, Z bosons
g gluon	meson
e electron	photon
m muon	star
t tau	galaxy
n neutrino	ion
	atom
	black hole

Particle Data Group, LBNL, © 2000. Supported by DOE and NSF

$T \sim 10^3$

n

At least 1 neutrino species is now NR



cosmic microwave radiation visible



n



n



n

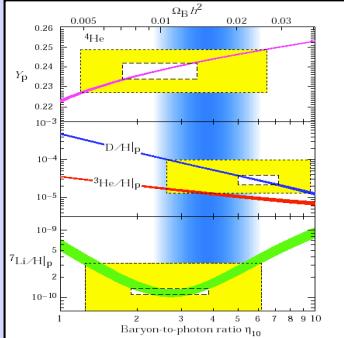
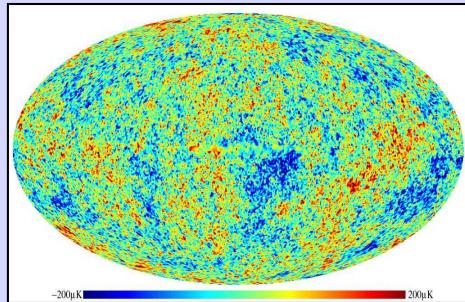
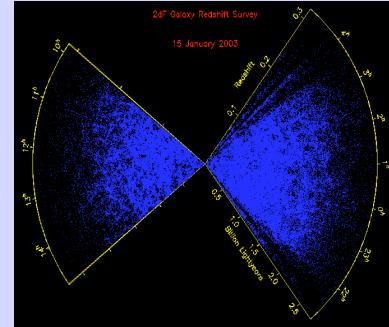


n



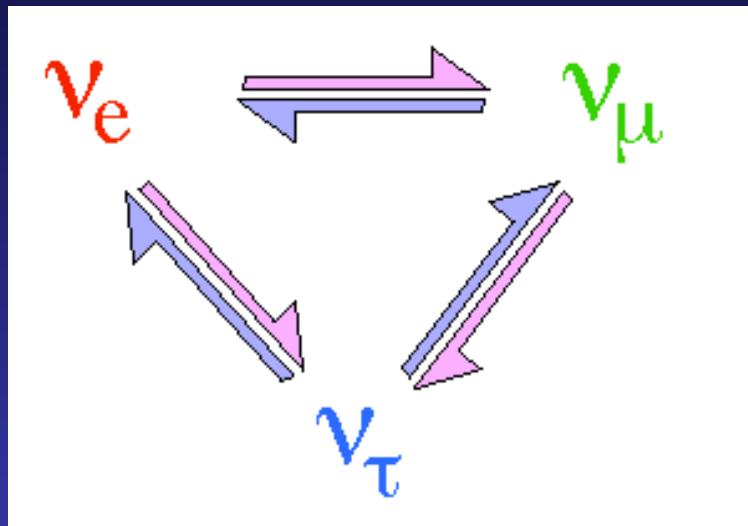
n

# Relic neutrinos influence several cosmological epochs

			
Primordial Nucleosynthesis BBN	Cosmic Microwave Background CMB	Formation of Large Scale Structures LSS	
$T \sim \text{MeV}$	$T < \text{eV}$		
$\nu_e$ vs $\nu_{\mu,\tau}$ $N_{\text{eff}}$	No flavour sensitivity		$N_{\text{eff}}$ & $m_\nu$

# We know that flavour neutrino oscillations exist

From present evidences of oscillations from experiments measuring atmospheric, solar, reactor and accelerator neutrinos



$$(e, \mu, \tau) \leftrightarrow (\nu_1, \nu_2, \nu_3)$$

$$|\nu_\alpha\rangle = \sum_i U_{\alpha i}^* |\nu_i\rangle$$

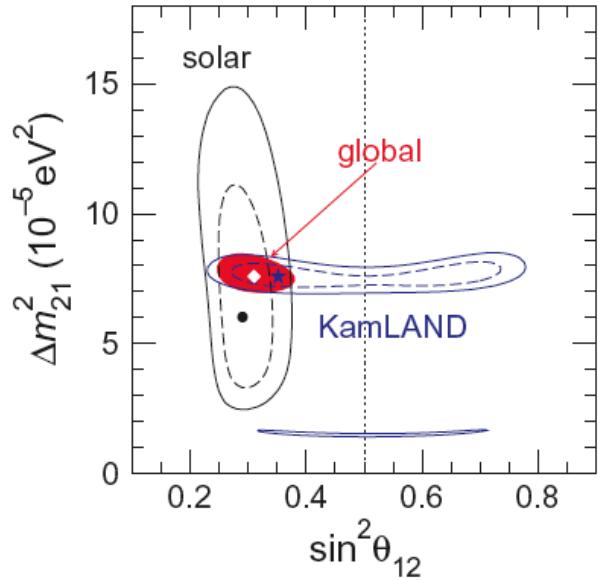
Evidence for Particle Physics  
beyond the Standard Model !

# Mixing Parameters...

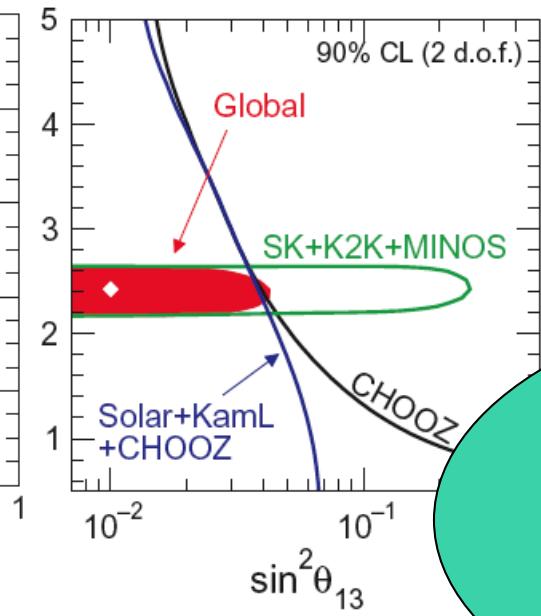
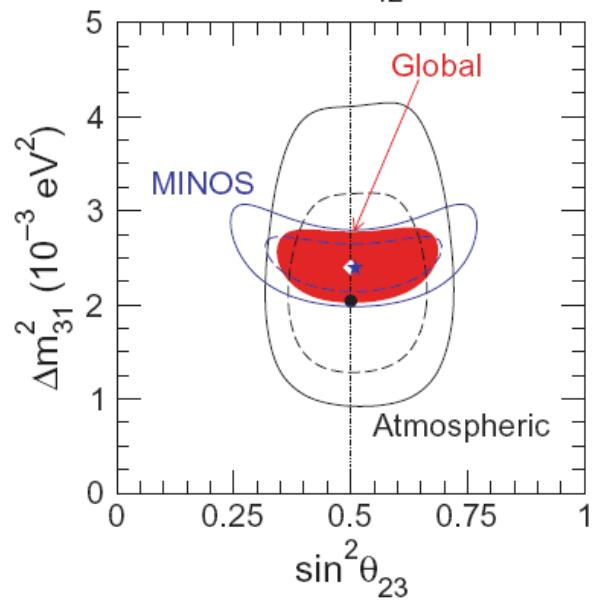
From present evidences of oscillations from experiments measuring atmospheric, solar, reactor and accelerator neutrinos

$$\begin{matrix} & \nu_1 & \nu_2 & \nu_3 \\ \nu_e & c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ \nu_\mu & -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ \nu_\tau & s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{matrix} \times \text{diag}(e^{i\alpha_1/2}, e^{i\alpha_2/2}, 1) .$$

# Mixing Parameters...



Parameter	Best fit	$2\sigma$	$3\sigma$
$\Delta m_{21}^2$ ( $10^{-5}$ eV $^2$ )	$7.65^{+0.23}_{-0.20}$	7.25–8.11	7.05–8.34
$ \Delta m_{31}^2 $ ( $10^{-3}$ eV $^2$ )	$2.40^{+0.12}_{-0.11}$	2.18–2.64	2.07–2.75
$\sin^2 \theta_{12}$	$0.304^{+0.022}_{-0.016}$	0.27–0.35	0.25–0.37
$\sin^2 \theta_{23}$	$0.50^{+0.07}_{-0.06}$	0.39–0.63	0.36–0.67
$\sin^2 \theta_{13}$	$0.01^{+0.016}_{-0.011}$	$\leq 0.040$	$\leq 0.056$



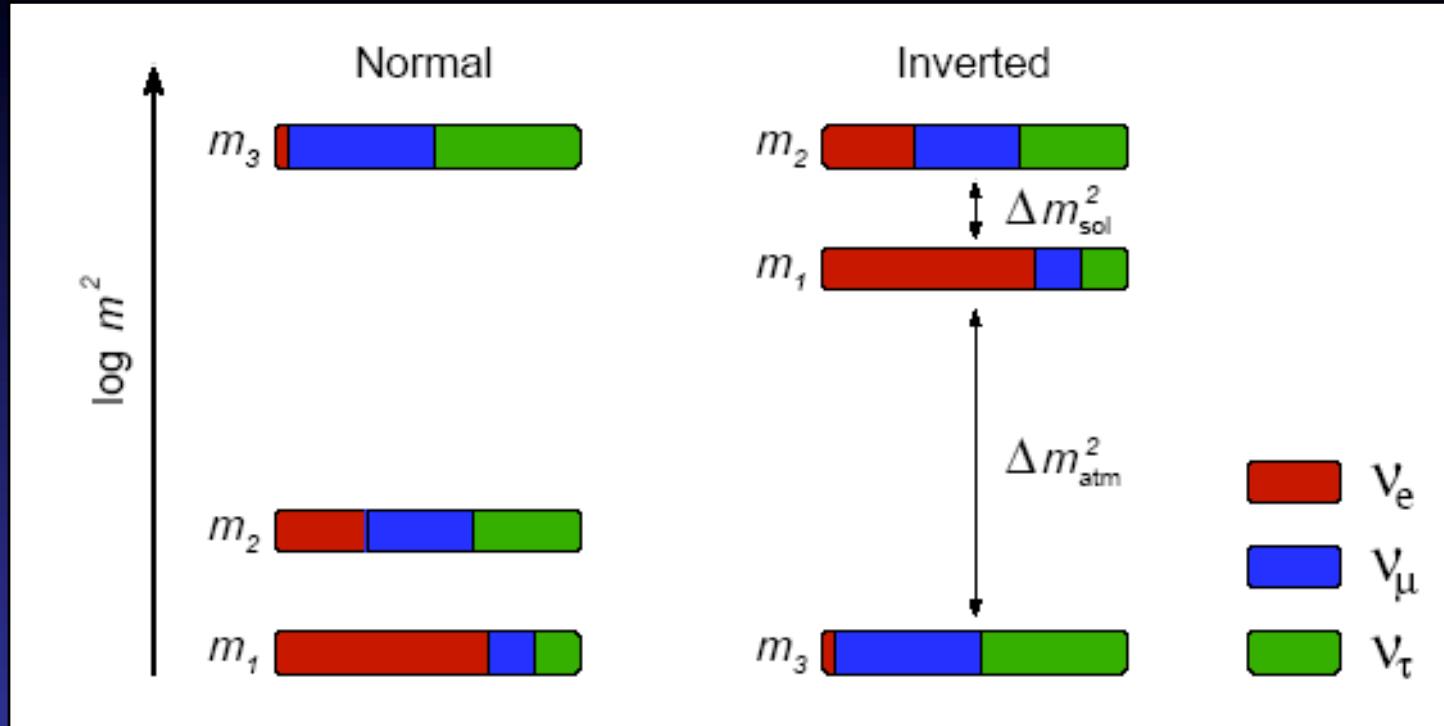
$$\nu_{\alpha L} = \sum_{i=1}^3 U_{\alpha i} \nu_{iL}, \quad (\alpha = e, \mu, \tau)$$

Present evidences  
for flavour neutrino  
oscillations: data on  
solar, atmospheric,  
reactor and accelerator  
neutrinos

Schwetz, Tórtola & Valle, New J Phys **10** (2008) 113011

[González-García & Maltoni , Phys Rep **460** (2008) 1,  
Fogli et al, Phys Rev Lett **101** (2008) 141801, and others]

## ... and neutrino masses



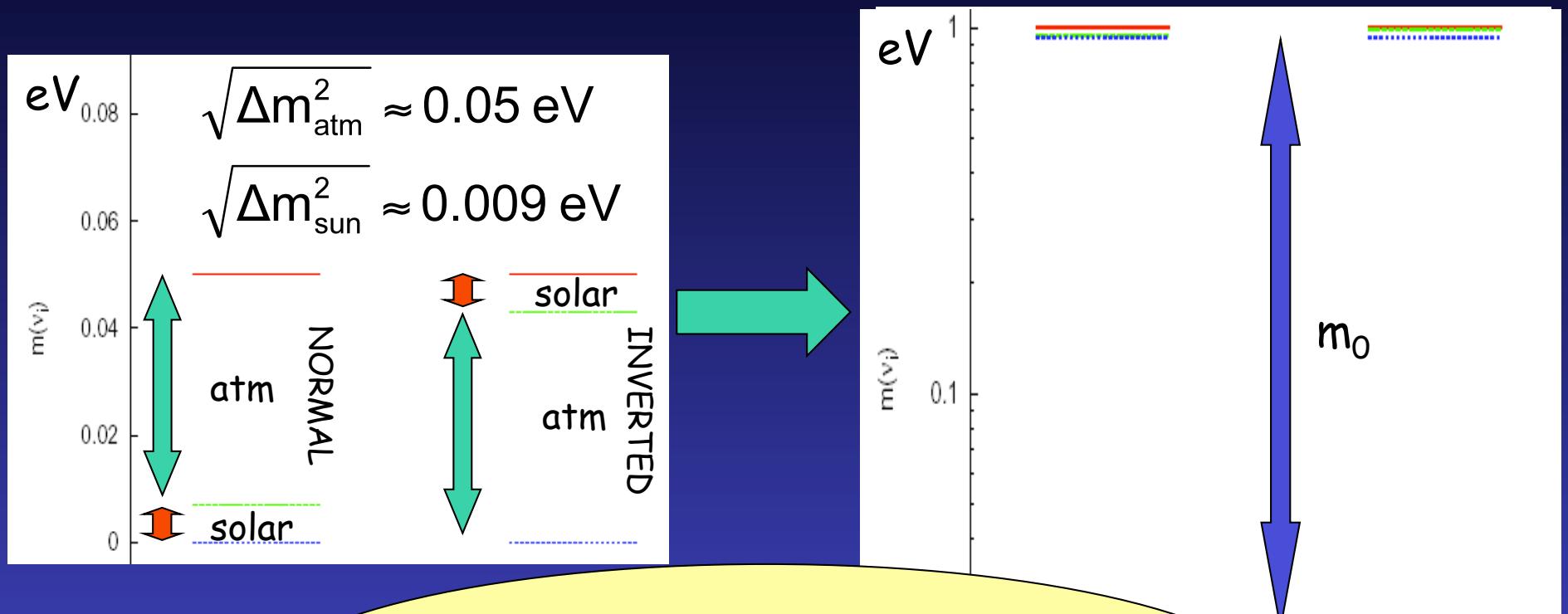
Possible neutrino mass hierarchy patterns

$$\nu_{\alpha L} = \sum_{i=1}^3 U_{\alpha i} \nu_{iL}, \quad (\alpha = e, \mu, \tau)$$

Present evidences  
for flavour neutrino  
oscillations: data on  
solar, atmospheric,  
reactor and accelerator  
neutrinos

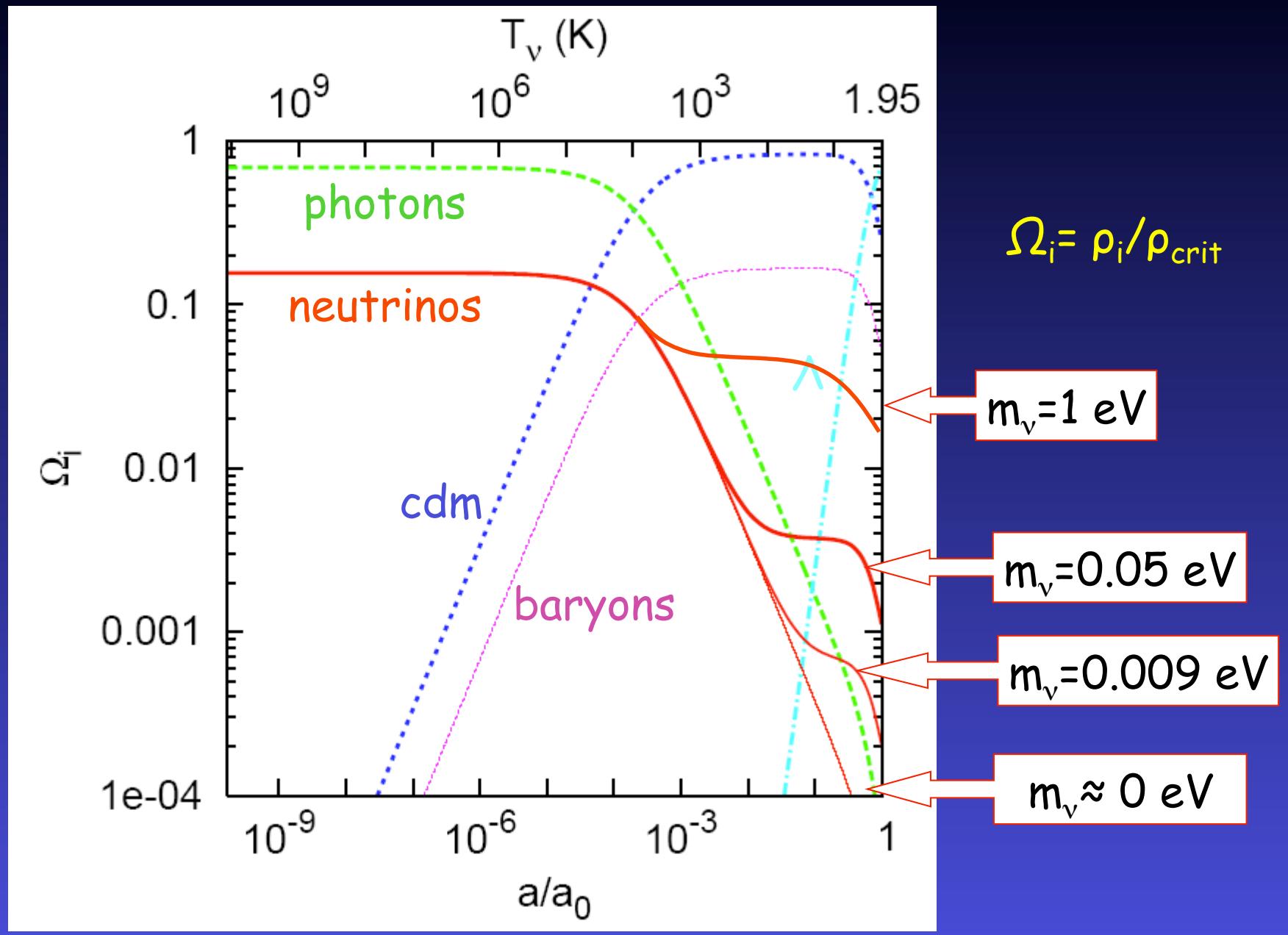
# ... and neutrino masses

Data on flavour oscillations do not fix the absolute scale of neutrino masses



What is the value of  $m_0$ ?

# Evolution of the background densities: 1 MeV → now



# The Cosmic Neutrino Background

Neutrinos decoupled at  $T \sim \text{MeV}$ , keeping a spectrum as that of a relativistic species

$$f_\nu(p, T) = \frac{1}{e^{p/T_\nu} + 1}$$

- Number density

At present  $112 (\nu + \bar{\nu}) \text{ cm}^{-3}$  per flavour

- Energy density

$$\Omega_\nu h^2 \simeq 1.7 \times 10^{-5} \quad \text{Massless}$$

Contribution to the energy density of the Universe

$$\Omega_\nu h^2 = \frac{\sum_i m_{\nu_i}}{94.1 \text{ eV}}$$

Massive  
 $m_\nu \gg T$

# Neutrinos as Dark Matter

- Neutrinos are natural **DM candidates**

$$\Omega_\nu h^2 = \frac{\sum m_i}{93.2 \text{ eV}} \quad \Omega_\nu < 1 \rightarrow \sum m_i < 46 \text{ eV}$$
$$\Omega_\nu < \Omega_m \approx 0.3 \rightarrow \sum m_i < 15 \text{ eV}$$

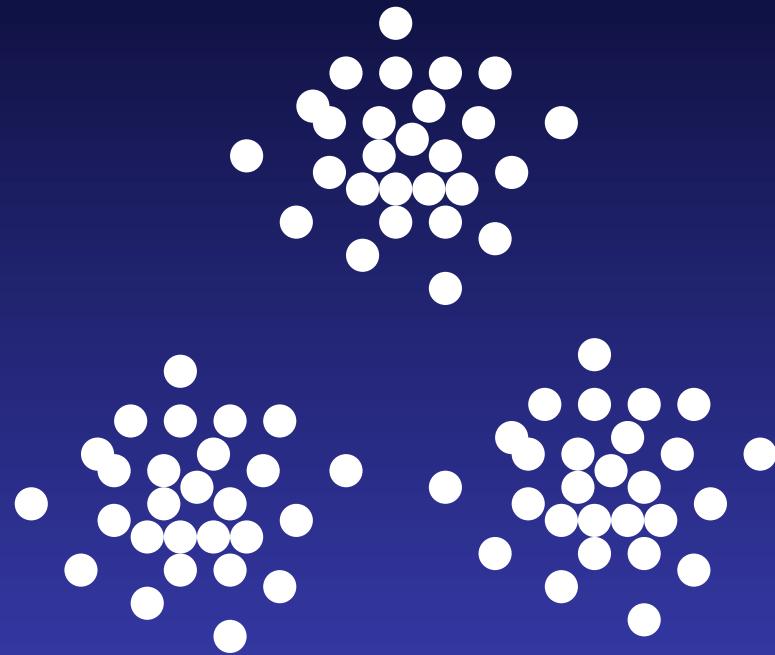
- They stream freely until non-relativistic (collisionless phase mixing)  **Neutrinos are HOT Dark Matter**
- First structures to be formed when Universe became matter-dominated are **very large**
- Ruled out by structure formation  **CDM**

# Neutrinos as Hot Dark Matter

Massive Neutrinos can still be subdominant DM: limits on  $m_\nu$  from Structure Formation (combined with other cosmological data)

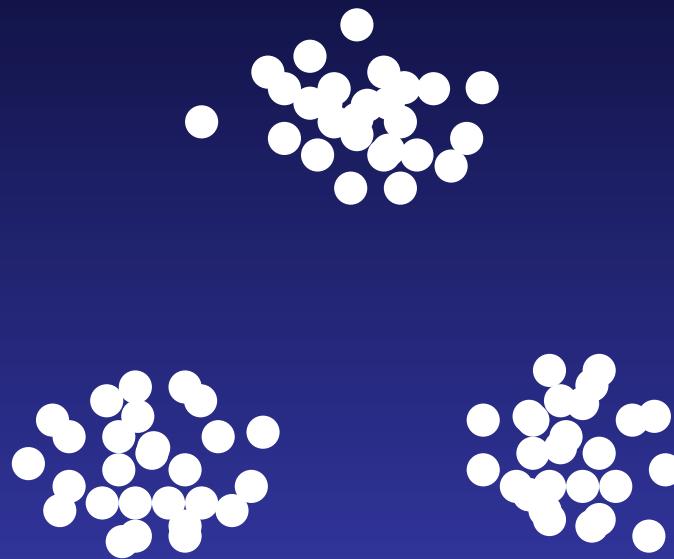
# Structure formation after equality

baryons and  
CDM (matter)  
experience  
gravitational  
clustering



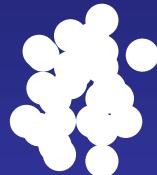
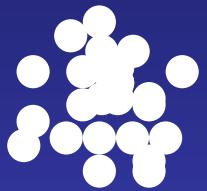
# Structure formation after equality

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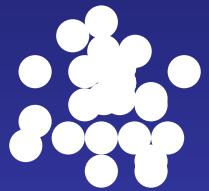
# Structure formation after equality

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# Structure formation after equality

baryons and  
CDM (matter)  
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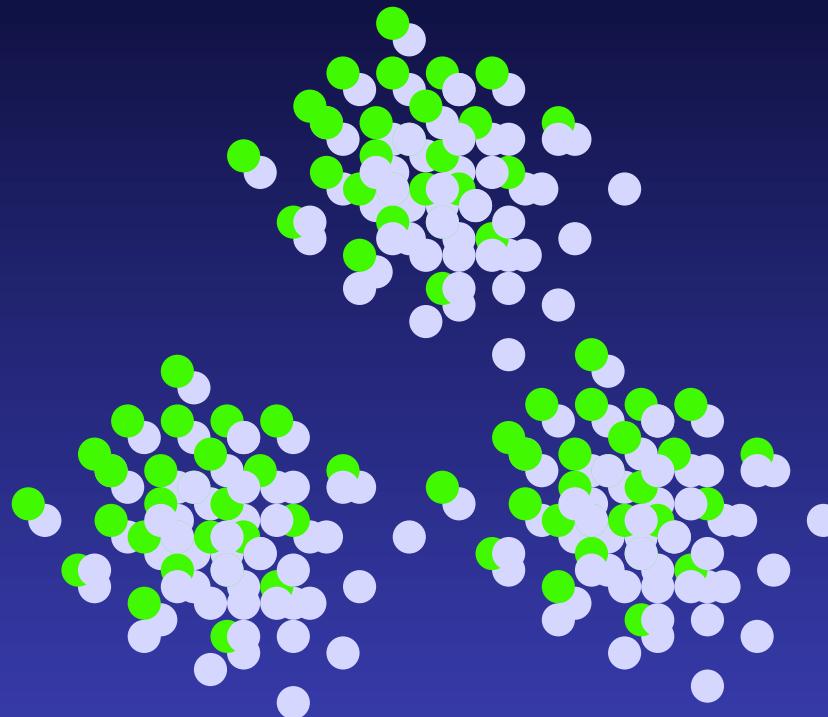


growth of  $\delta\rho/\rho(k,t)$  fixed by  
**gravity vs expansion** balance

$$\Rightarrow \delta\rho/\rho \propto a$$

# Structure formation after equality

baryons and  
CDM (matter)  
experience  
gravitational  
clustering

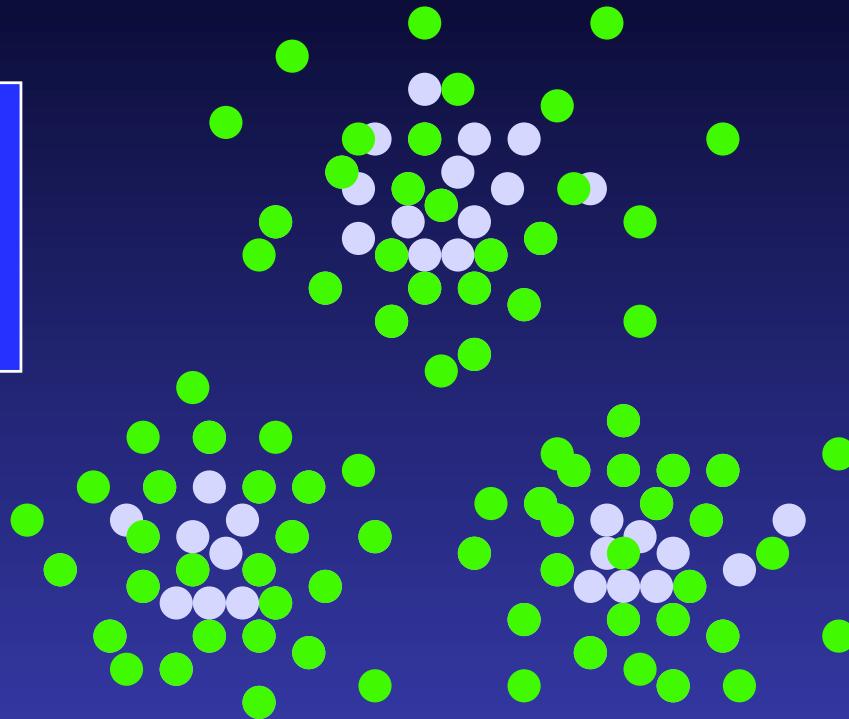


neutrinos  
experience  
free-streaming  
with  
 $v = c$  or  $\langle p \rangle/m$

# Structure formation after equality

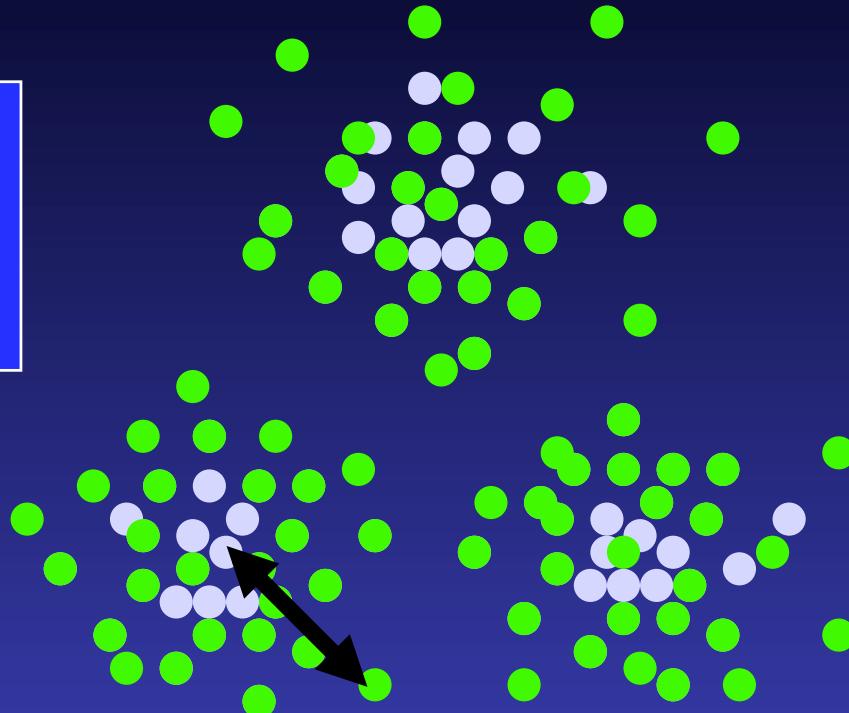
baryons and  
CDM (matter)  
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clustering

neutrinos  
experience  
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 $v = c$  or  $\langle p \rangle/m$



# Structure formation after equality

baryons and  
CDM (matter)  
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clustering



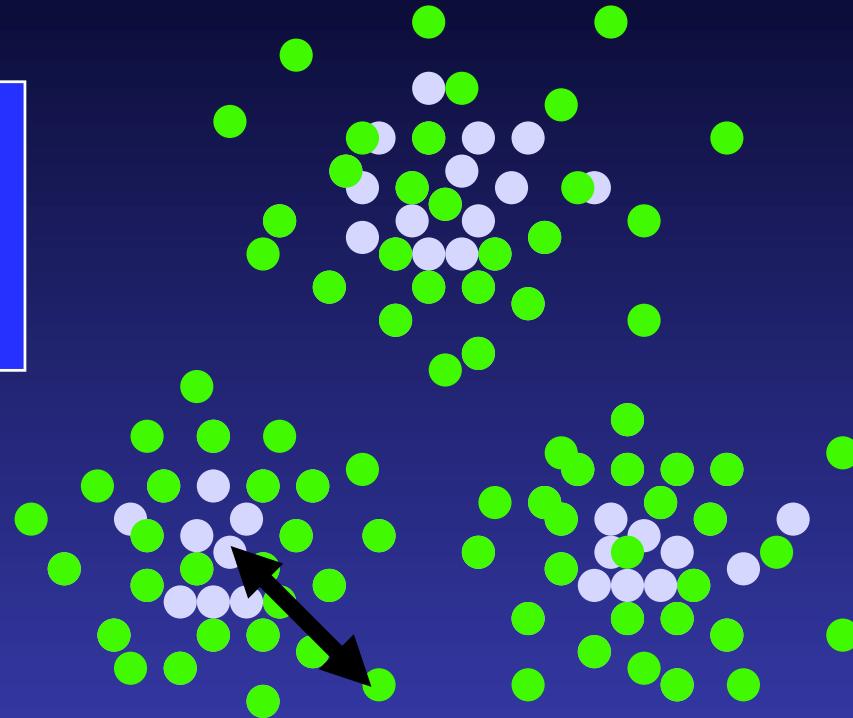
neutrinos  
experience  
free-streaming  
with  
 $v = c$  or  $\langle p \rangle/m$

neutrinos cannot cluster below a diffusion length

$$\lambda = \int v dt < \int c dt$$

# Structure formation after equality

baryons and  
CDM (matter)  
experience  
gravitational  
clustering



neutrinos  
experience  
free-streaming  
with  
 $v = c$  or  $\langle p \rangle / m$

$$\text{for } (2\pi/k) < \lambda ,$$

free-streaming suppresses growth of structures during MD

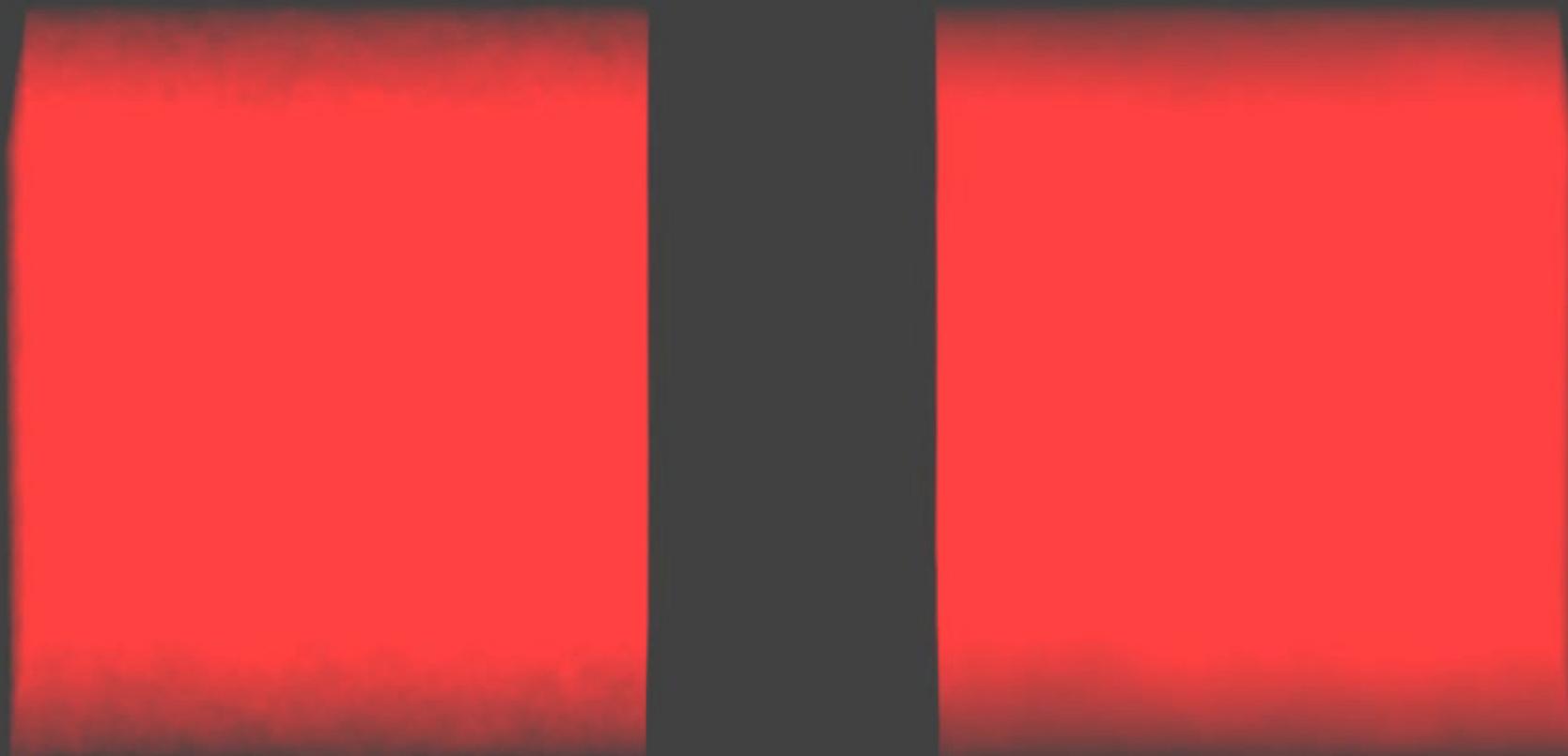
$$\Rightarrow \delta\rho/\rho \propto a^{1-3/5} f_v$$

$$\text{with } f_v = \rho_v / \rho_m \approx (\Sigma m_v) / (15 \text{ eV})$$

# Neutrinos as Hot Dark Matter

Massive Neutrinos can still be subdominant DM: limits on  $m_\nu$  from Structure Formation (combined with other cosmological data)

$Z=32.33$



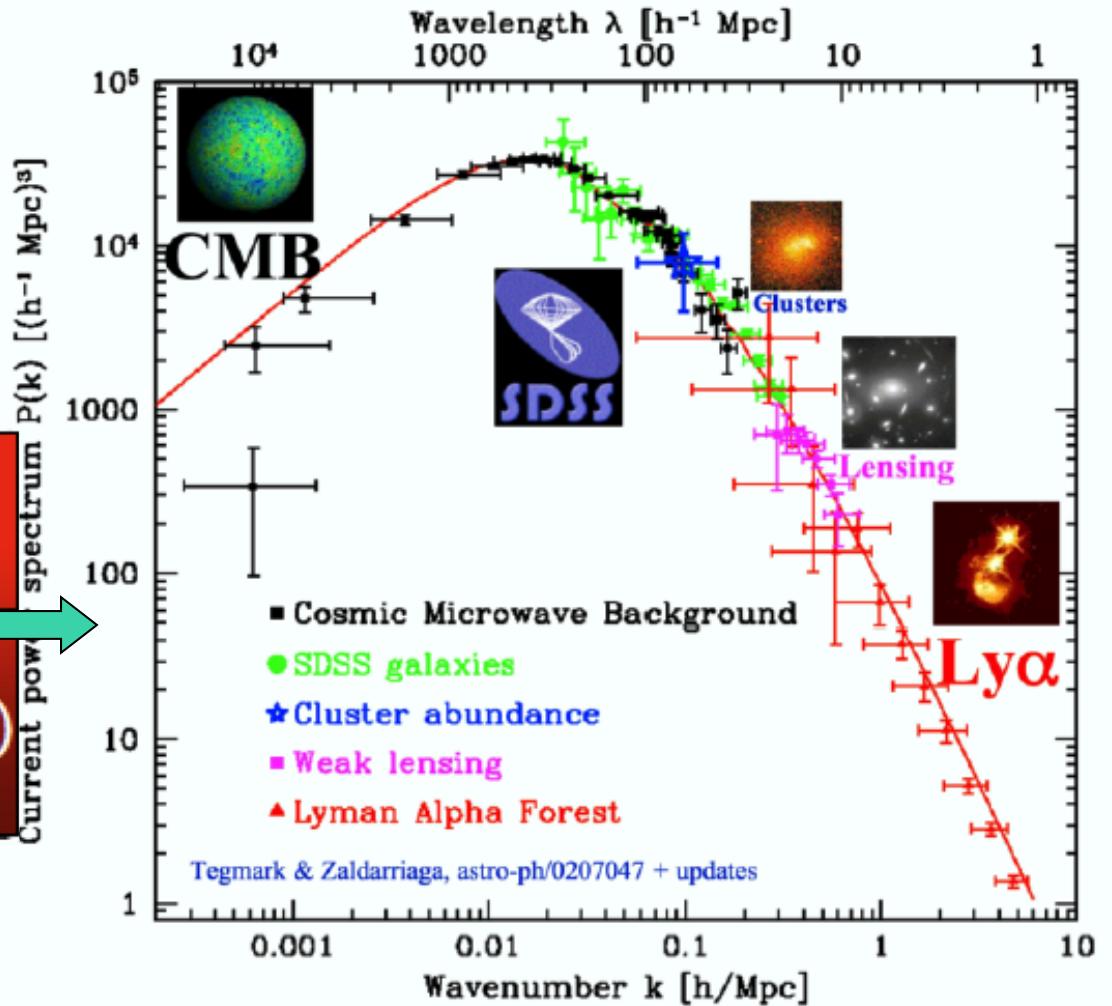
# Power Spectrum of density fluctuations

Field of density Fluctuations

$$\delta(x) = \frac{\delta\rho(x)}{\bar{\rho}}$$

Matter power spectrum is the Fourier transform of the two-point correlation function

$$\langle \delta(x_1)\delta(x_2) \rangle = \int \frac{d^3k}{(2\pi)^3} e^{ik(x_2-x_1)} P(k)$$



# Neutrinos as Hot Dark Matter: effect on $P(k)$

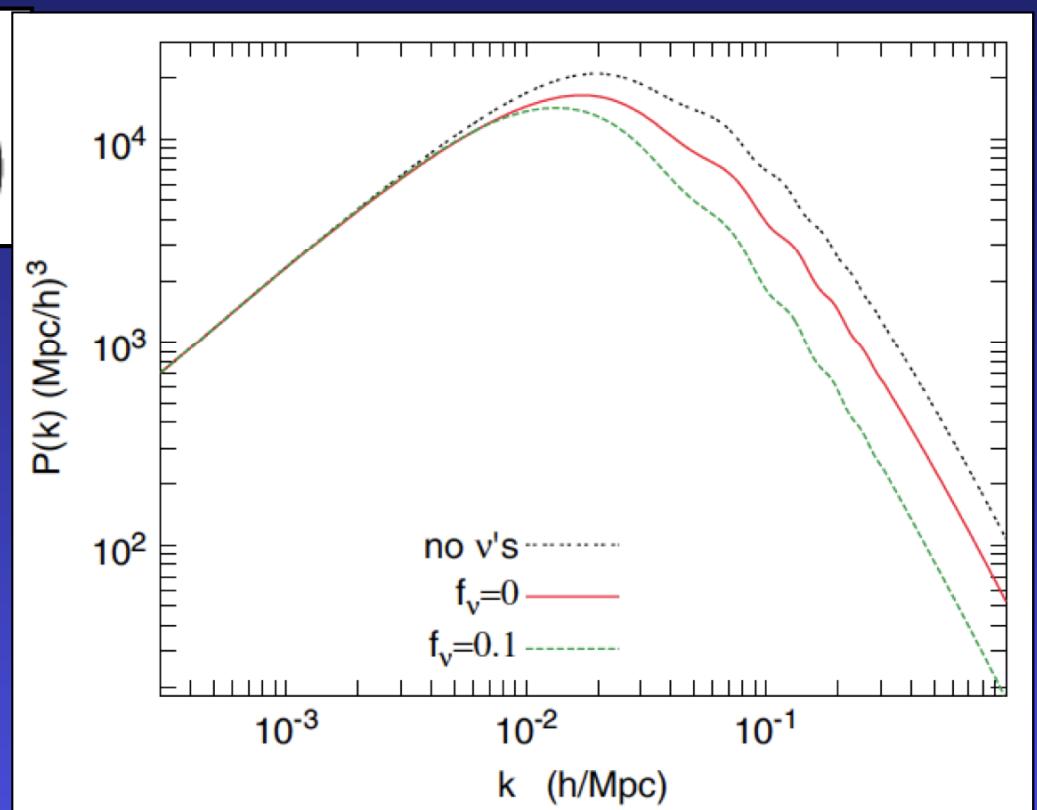
Massive Neutrinos can still be subdominant DM: limits on  $m_\nu$  from Structure Formation (combined with other cosmological data)

- Effect of Massive Neutrinos: suppression of Power at small scales

The small-scale suppression is given by

$$\left(\frac{\Delta P}{P}\right) \approx -8 \frac{\Omega_\nu}{\Omega_m} \approx -0.8 \left(\frac{m_\nu}{1 \text{ eV}}\right) \left(\frac{0.1N}{\Omega_m h^2}\right)$$

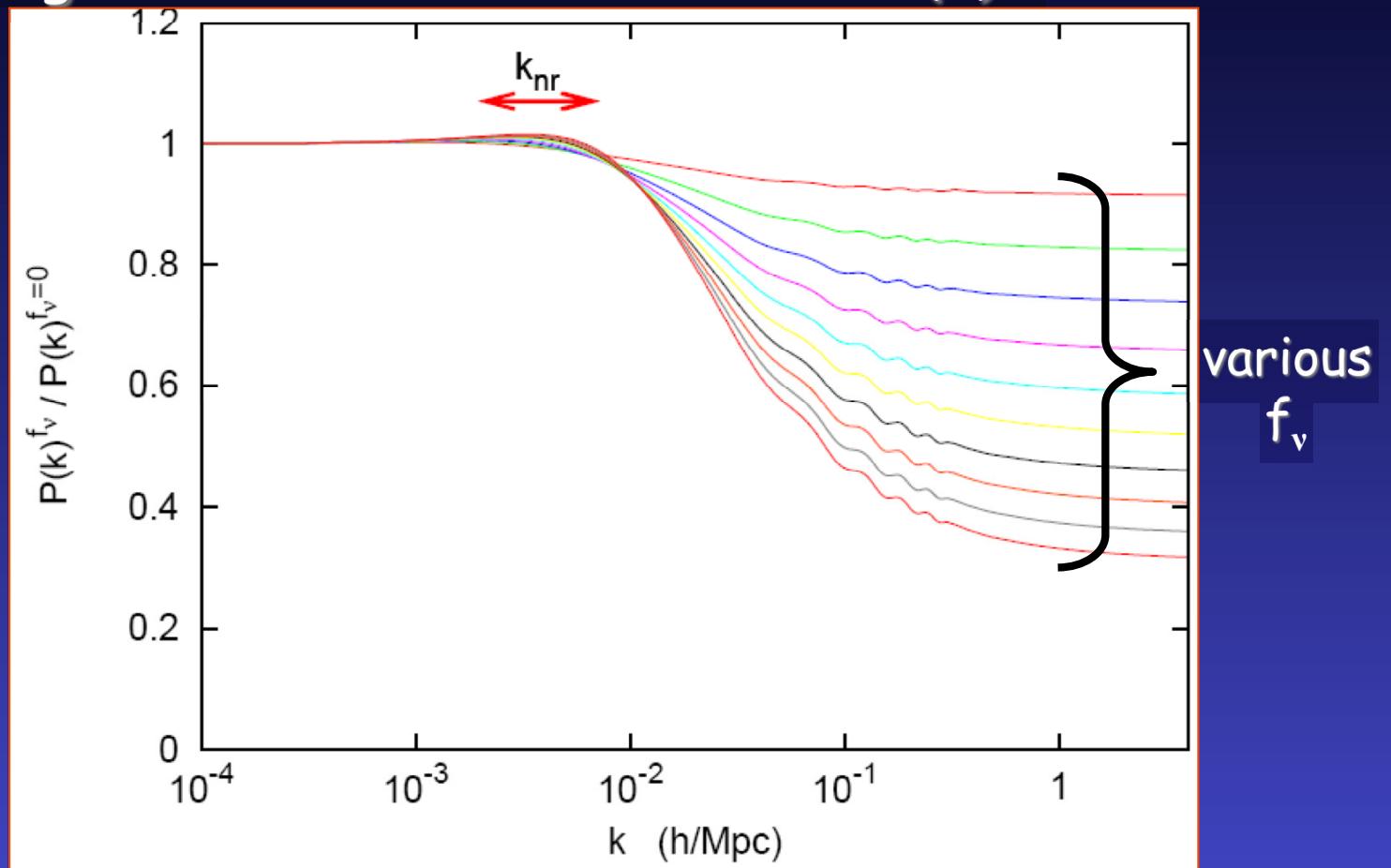
$f_\nu$



# Effect of massive neutrinos on $P(k)$

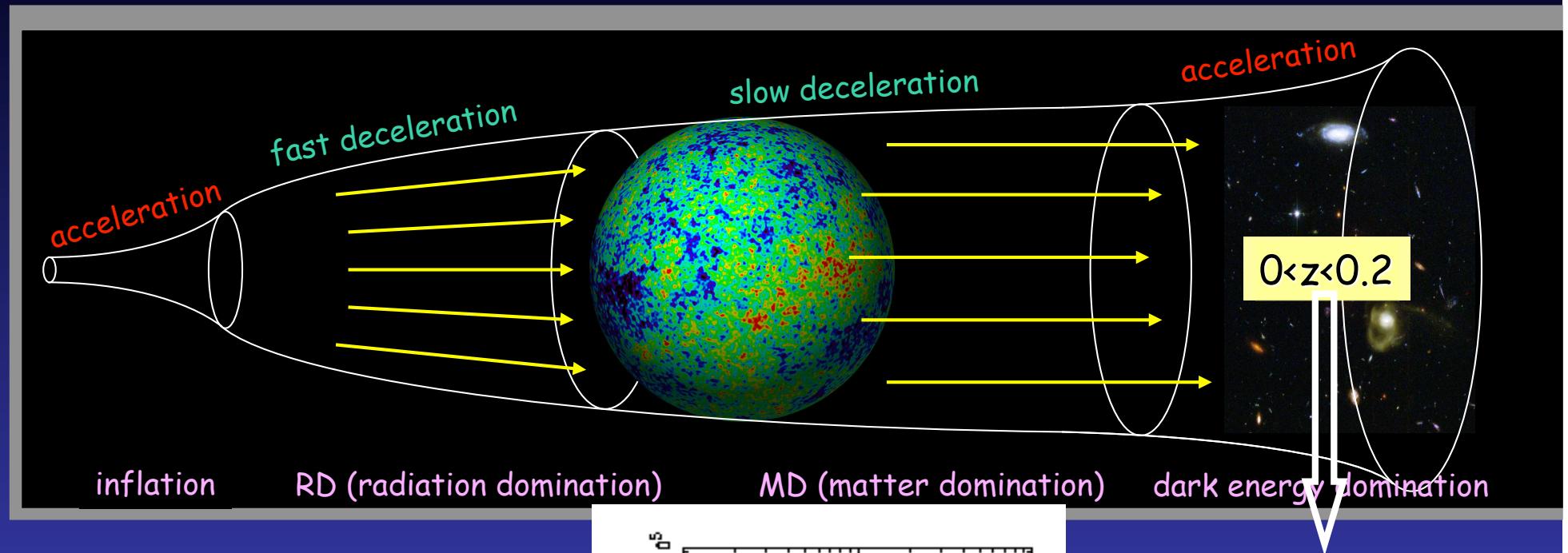
Observable signature of the total mass on  $P(k)$  :

$$\frac{P(k) \text{ massive}}{P(k) \text{ massless}}$$

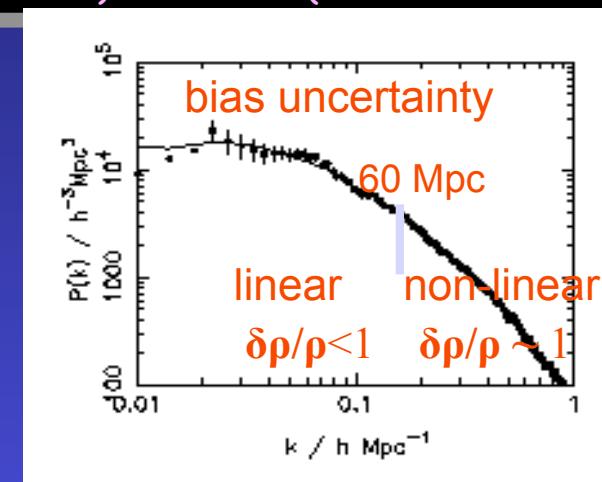


Lesgourgues & SP, Phys. Rep. 429 (2006) 307

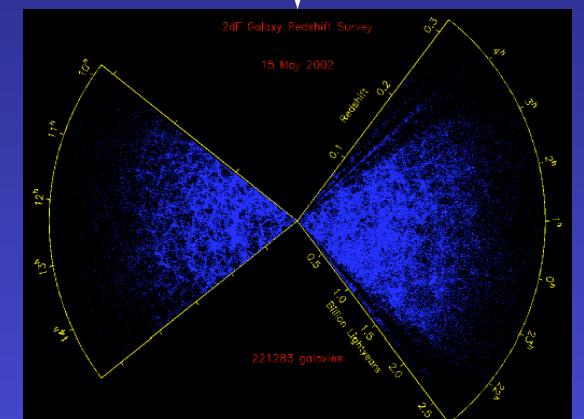
# Cosmological observables: LSS



Distribution  
of large-scale  
structures at low  $z$

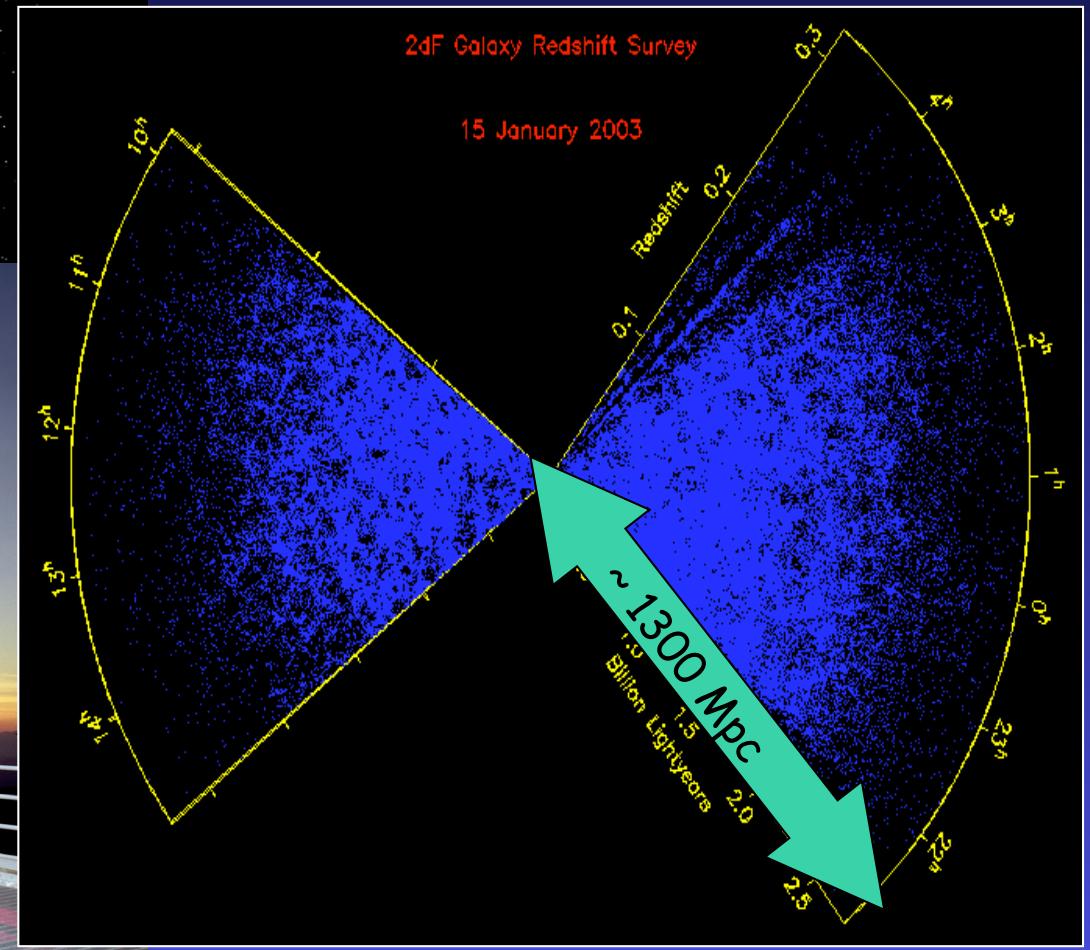


matter power spectrum  $P(k)$

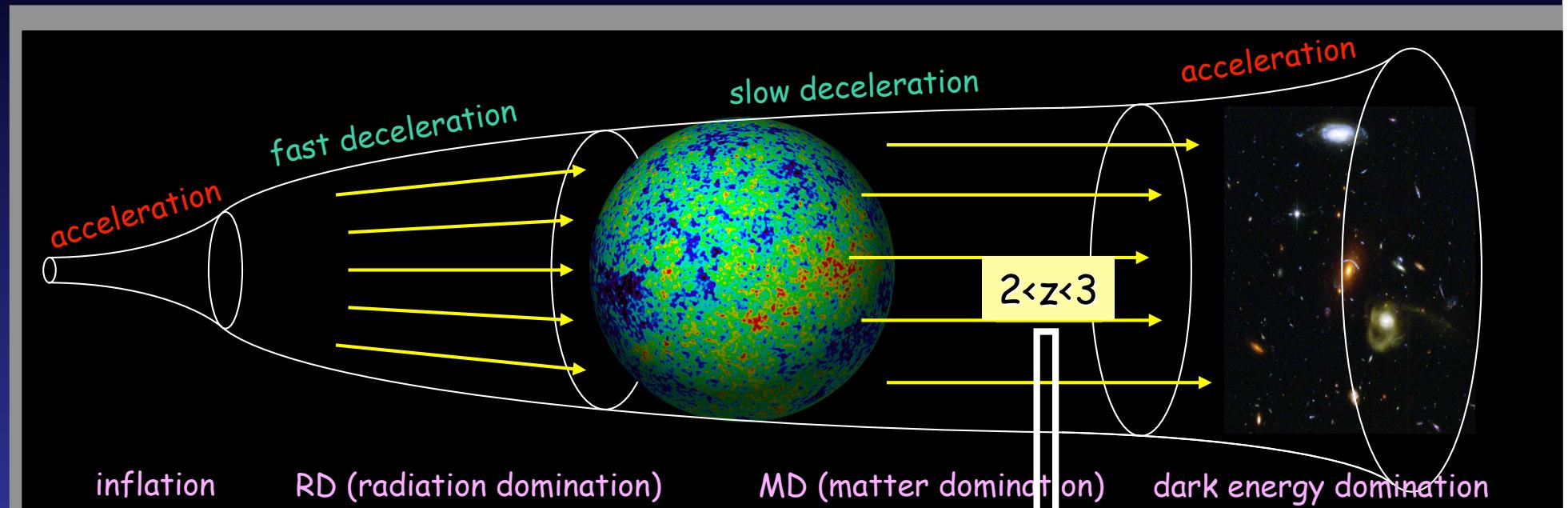


galaxy redshift surveys

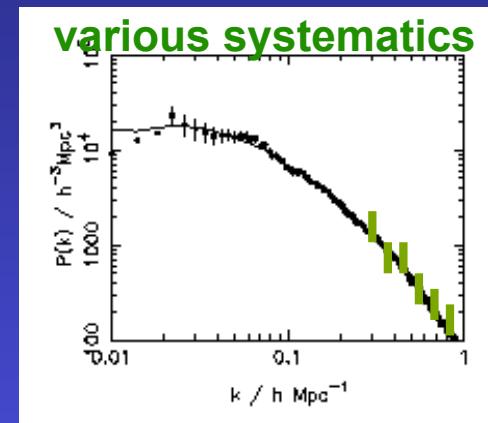
# Galaxy Redshift Surveys



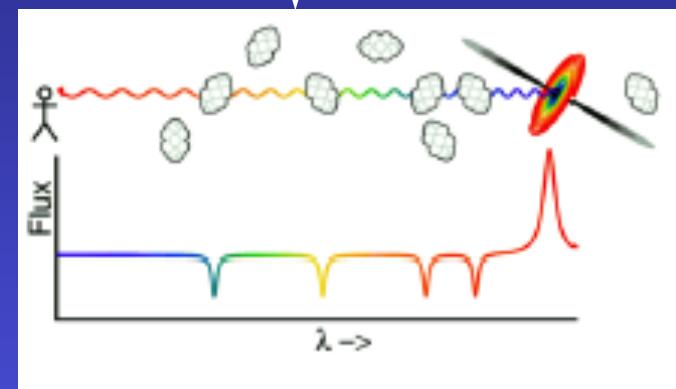
# Cosmological observables : LSS



Distribution  
of large-scale  
structures at  
medium  $z$

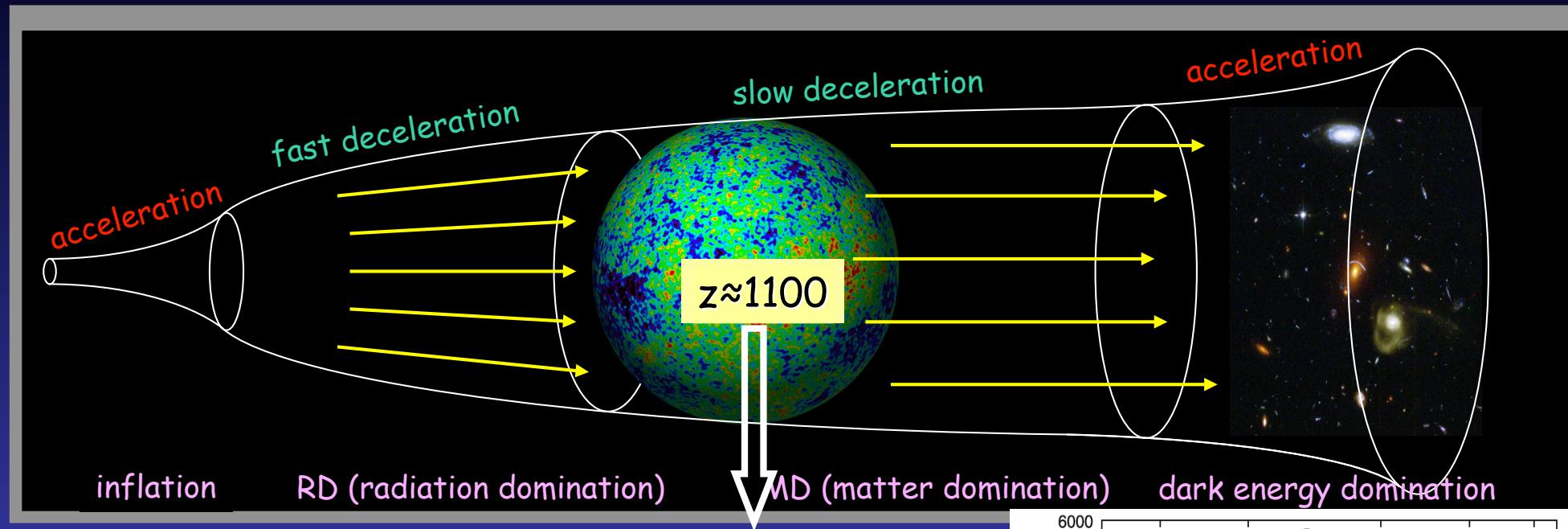


matter power spectrum  $P(k)$

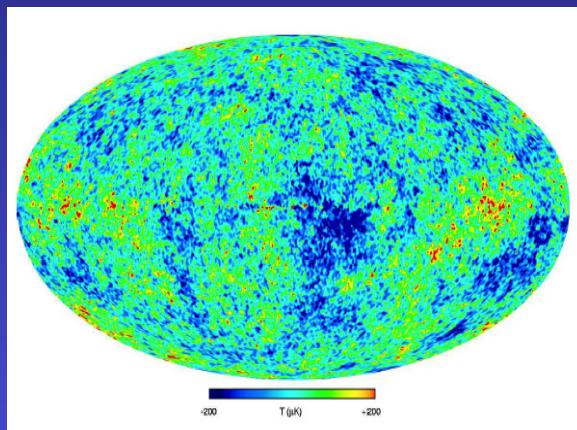


Lyman- $\alpha$  forests in quasar spectra

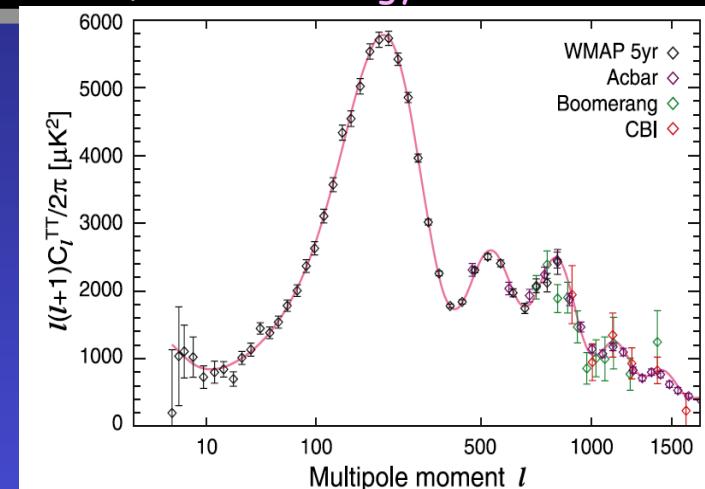
# Cosmological observables: CMB



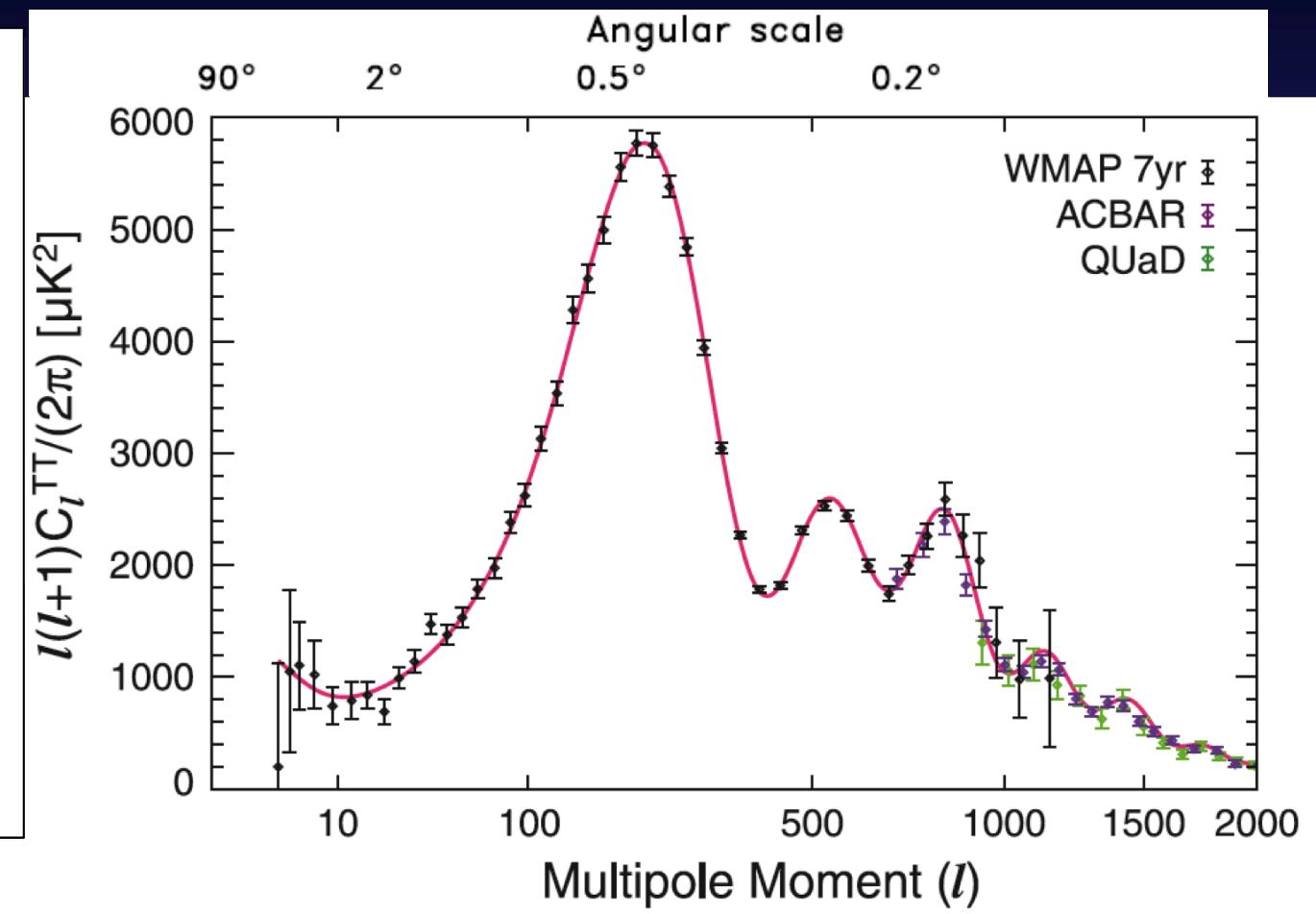
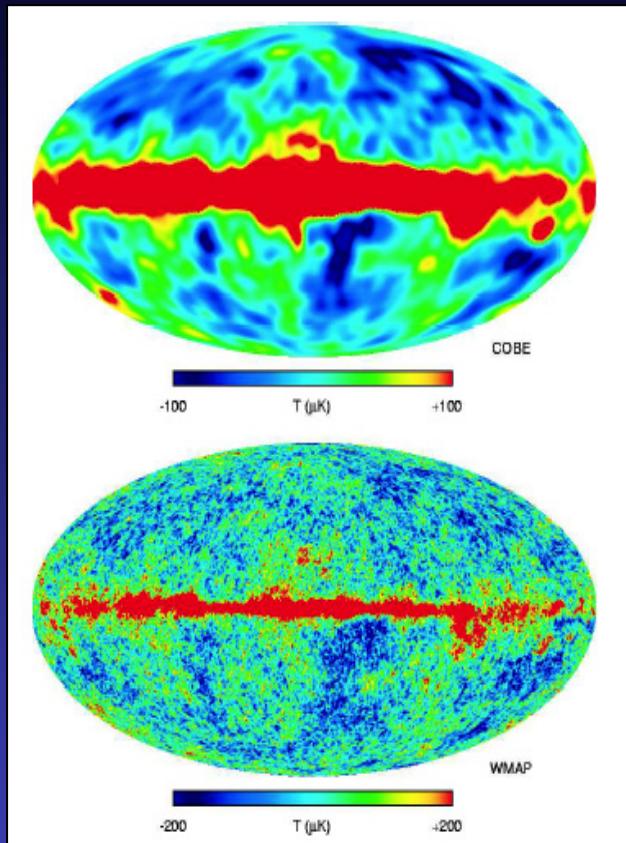
Anisotropies  
of the Cosmic  
Microwave  
Background



CMB temperature/polarization anisotropies  $\Rightarrow$  photon power spectra



# CMB TT DATA



# Effect of massive neutrinos on the CMB spectra

- 1) CMB spectrum essentially unchanged if neutrinos become NR AFTER photon decoupling ( $z_{\text{rec}} \sim 1089$ )

$$\begin{aligned} 1 + z_{\text{nr}} &= \frac{T_{\nu, \text{nr}}}{T_{\nu, 0}} \\ &= 1.99 \times 10^3 (m_\nu / \text{eV}) \\ &= 6.24 \times 10^4 \omega_\nu, \end{aligned}$$

Neutrinos become NR BEFORE recombination if:

$$\omega_\nu \gtrsim 0.017 \implies \sum_i m_{\nu_i} \gtrsim 1.6 \text{ eV}$$

More details including effects of neutrino mass on “reduced CMB observables” in Ichikawa et al, PRD 71 (2005) 043001

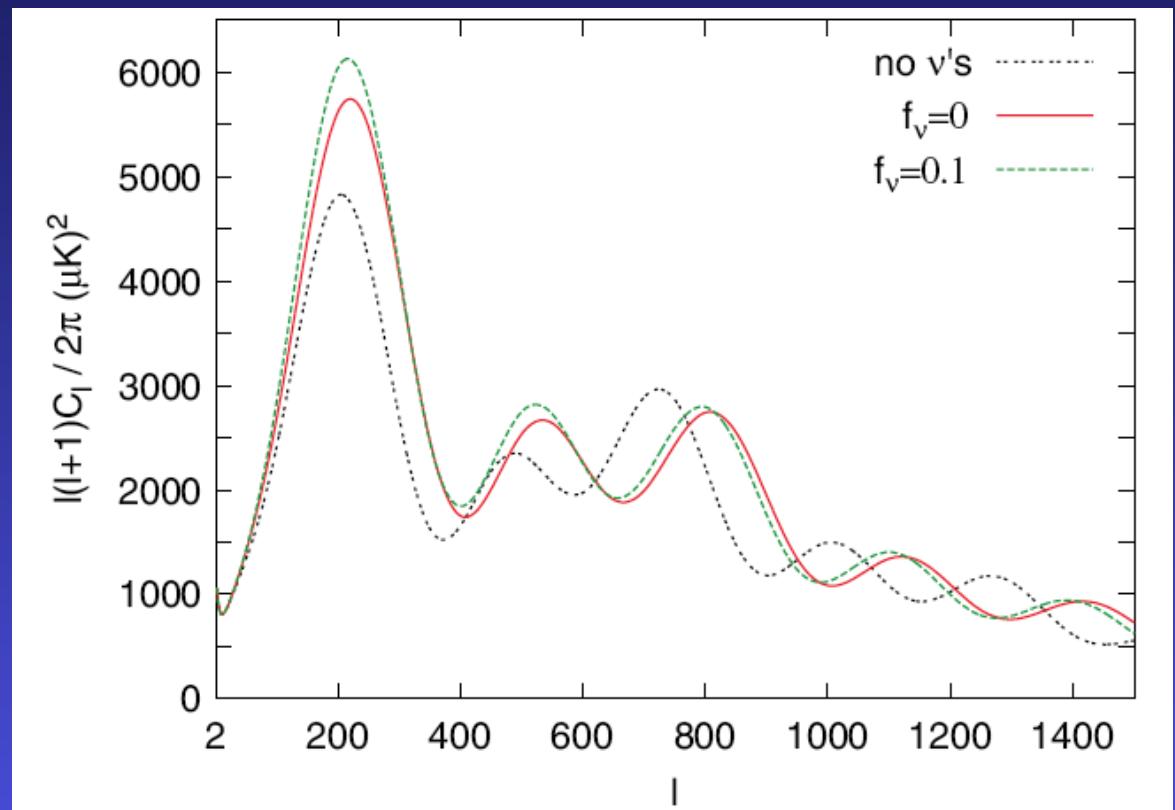
# Effect of massive neutrinos on the CMB spectra

- 1) CMB spectrum essentially unchanged if neutrinos become NR AFTER photon decoupling.
- 2) Impact on CMB spectra is indirect: non-zero  $\Omega_\nu$  modifies the background evolution (change in equality time)

Ex: in a flat universe,

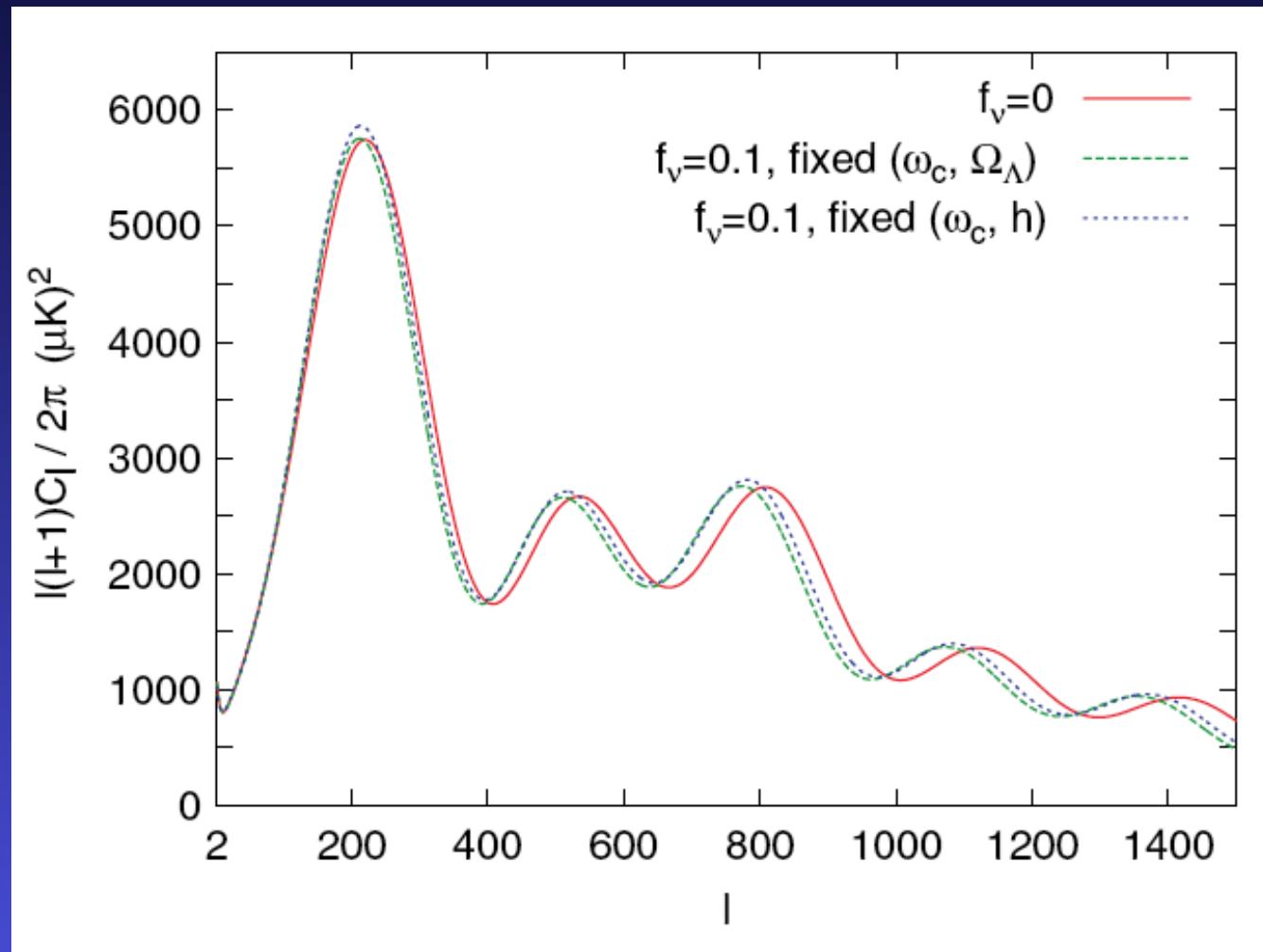
keep  $\Omega_\Lambda + \Omega_{\text{cdm}} + \Omega_b + \Omega_\nu = 1$

constant



# Effect of massive neutrinos on the CMB spectra

Problem with **parameter degeneracies**: change in other cosmological parameters can mimic the effect of nu masses



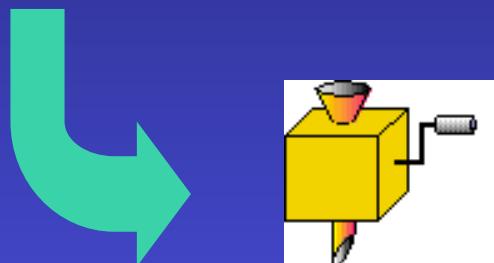
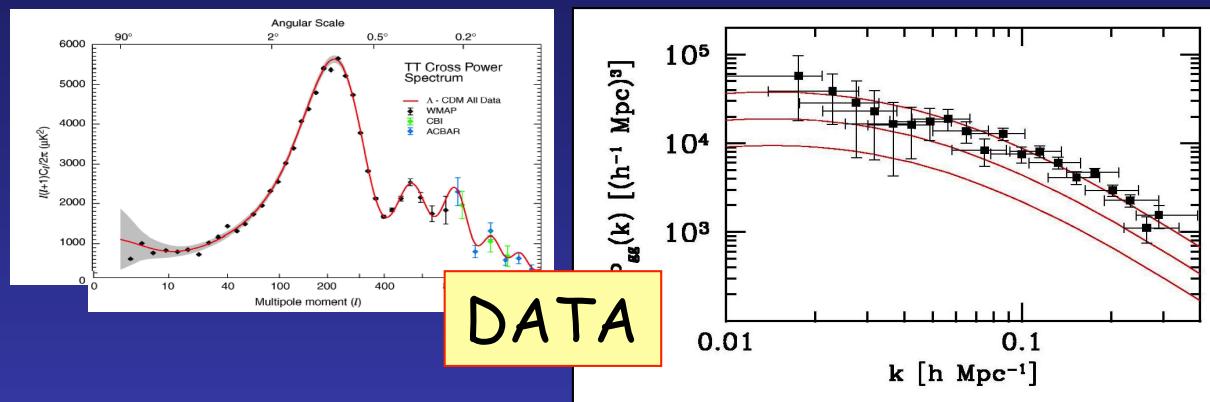
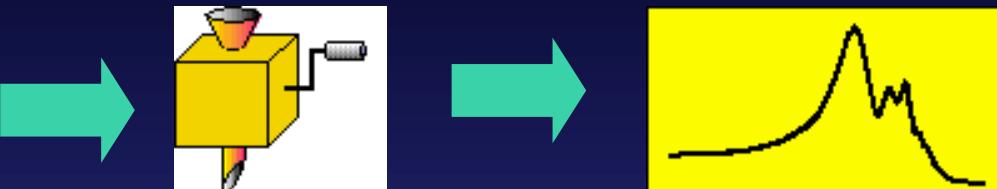
**Bounds on  $m_\nu$  from CMB,  
LSS and other data**

# Neutrinos as Hot Dark Matter

Massive Neutrinos can still be subdominant DM: limits on  $m_\nu$  from Structure Formation (combined with other cosmological data)

# How to get a bound (measurement) of neutrino masses from Cosmology

Fiducial cosmological model:  
 $(\Omega_b h^2, \Omega_m h^2, h, n_s, \tau, \Sigma m_\nu)$



PARAMETER  
ESTIMATES

# Cosmological Data

- CMB Temperature: WMAP plus data from other experiments at large multipoles (CBI, ACBAR, VSA...)
- CMB Polarization: WMAP, ...
- Large Scale Structure:
  - \* Galaxy Clustering (2dF, SDSS)
  - \* Bias (Galaxy, ...): Amplitude of the Matter  $P(k)$  ( $SDSS, \sigma_8$ )
  - \* Lyman-a forest: independent measurement of power on small scales
  - \* Baryon acoustic oscillations (SDSS)

Bounds on parameters from other data: SNIa ( $\Omega_m$ ), HST ( $h$ ), ...

# Cosmological Parameters: example

Parameter	Meaning	Status
$\tau$	Reionization optical depth	Not optional
$\omega_b$	Baryon density	Not optional
$\omega_d$	Dark matter density	Not optional
$f_\nu$	Dark matter neutrino fraction	Well motivated
$\Omega_\Lambda$	Dark energy density	Not optional
$w$	Dark energy equation of state	Worth testing
$\Omega_k$	Spatial curvature	Worth testing
$A_s$	Scalar fluctuation amplitude	Not optional
$n_s$	Scalar spectral index	Well motivated
$\alpha$	Running of spectral index	Worth testing
$r$	Tensor-to-scalar ratio	Well motivated
$n_t$	Tensor spectral index	Well motivated
$b$	Galaxy bias factor	Not optional

SDSS Coll, PRD 69 (2004) 103501

## Cosmological bounds on neutrino mass(es)

A unique cosmological bound on  $m_\nu$  DOES NOT exist !

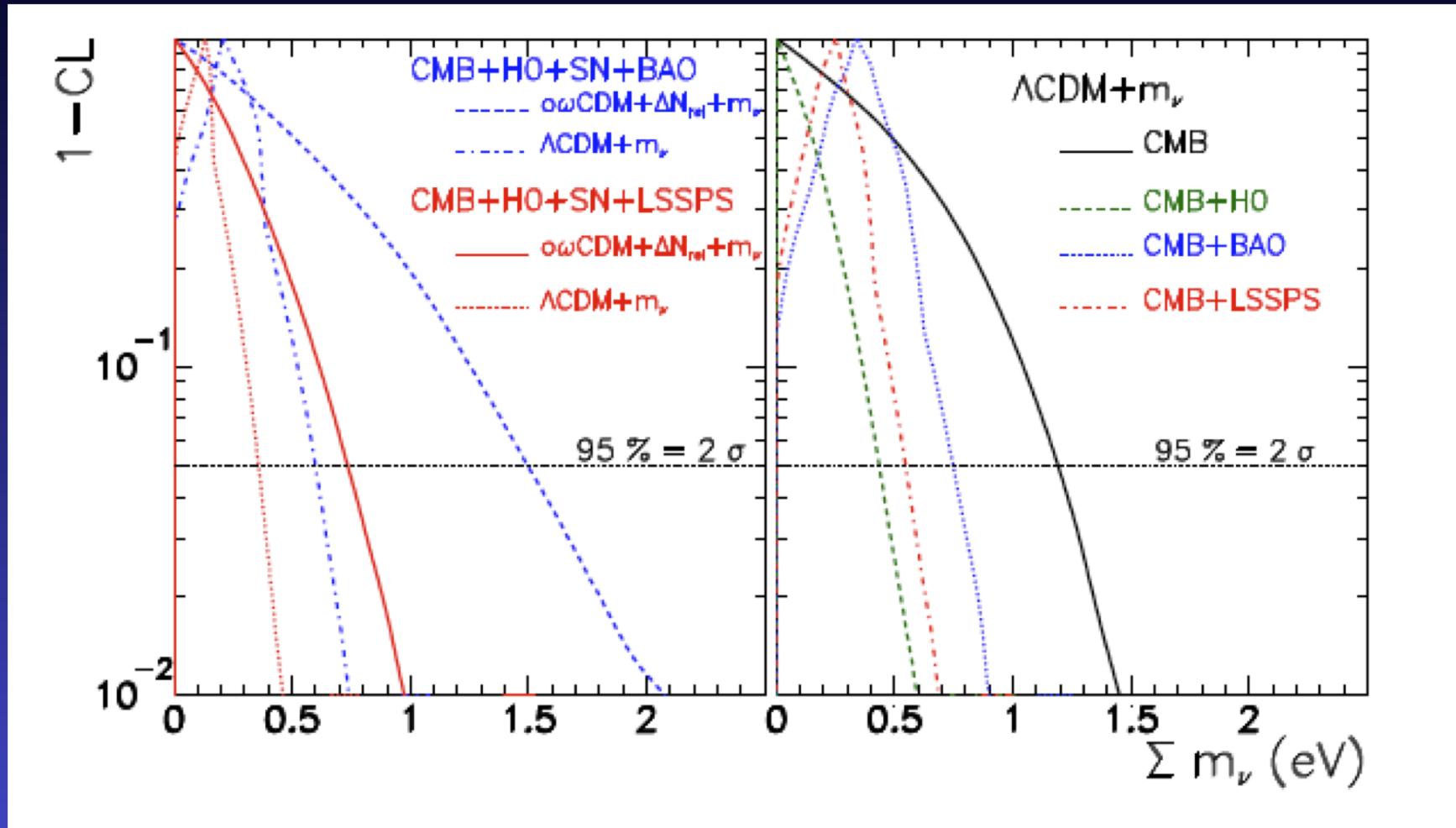
# Cosmological bounds on neutrino mass(es)

A unique cosmological bound on  $m_\nu$  DOES NOT exist !

Different analyses have found upper bounds on neutrino masses, since they depend on

- The combination of cosmological data used
- The assumed cosmological model: number of parameters (problem of parameter degeneracies)
- The properties of relic neutrinos

## Current cosmological bounds on neutrino masses



# Current cosmological bounds on neutrino masses

	CMB+HO+SN+BAO			CMB+HO+SN+LSS-PS		
	best	$1\sigma$	95% CL	best	$1\sigma$	95% CL
$H_0$ km/s/Mpc	76.2	$+3.0$ $-2.8$	$+5.7$ $-5.6$	74.4	$+2.8$ $-2.9$	$+5.6$ $-5.6$
$\Omega_b h^2 \times 100$	2.205	$+0.057$ $-0.050$	$+0.103$ $-0.105$	2.239	$+0.059$ $-0.046$	$+0.095$ $-0.108$
$\Omega_c h^2$	0.131	$+0.018$ $-0.013$	$+0.036$ $-0.023$	0.128	$+0.024$ $-0.009$	$+0.042$ $-0.018$
$n_S$	0.961	$+0.021$ $-0.015$	$+0.040$ $-0.030$	0.971	$+0.019$ $-0.017$	$+0.037$ $-0.033$
$\tau$	0.086	$+0.011$ $-0.015$	$+0.026$ $-0.028$	0.083	$+0.016$ $-0.011$	$+0.030$ $-0.023$
$\sigma_8$	0.787	$+0.091$ $-0.073$	$+0.135$ $-0.179$	0.824	$+0.051$ $-0.048$	$+0.097$ $-0.105$
$\Omega_k$	-0.006	$+0.010$ $-0.009$	$-0.022 \leq \Omega_k \leq 0.016$	-0.011	$+0.008$ $-0.009$	$-0.028 \leq \Omega_k \leq 0.007$
$\omega$	-1.17	$+0.19$ $-0.21$	$-0.62 \leq \omega + 1 \leq 0.18$	-1.12	$+0.21$ $-0.20$	$-0.57 \leq \omega + 1 \leq 0.26$
$\Delta N_{\text{rel}}$	1.2	$+1.1$ $-0.61$	$0.08 \leq \Delta N_{\text{rel}} \leq 3.2$	1.3	$+1.4$ $-0.54$	$0.21 \leq \Delta N_{\text{rel}} \leq 3.6$
$\sum m_\nu$ (eV)		$\leq 0.77$	$\leq 1.5$		$\leq 0.37$	$\leq 0.76$

# Cosmological bounds on neutrino masses using WMAP

Dependence on the data set used:

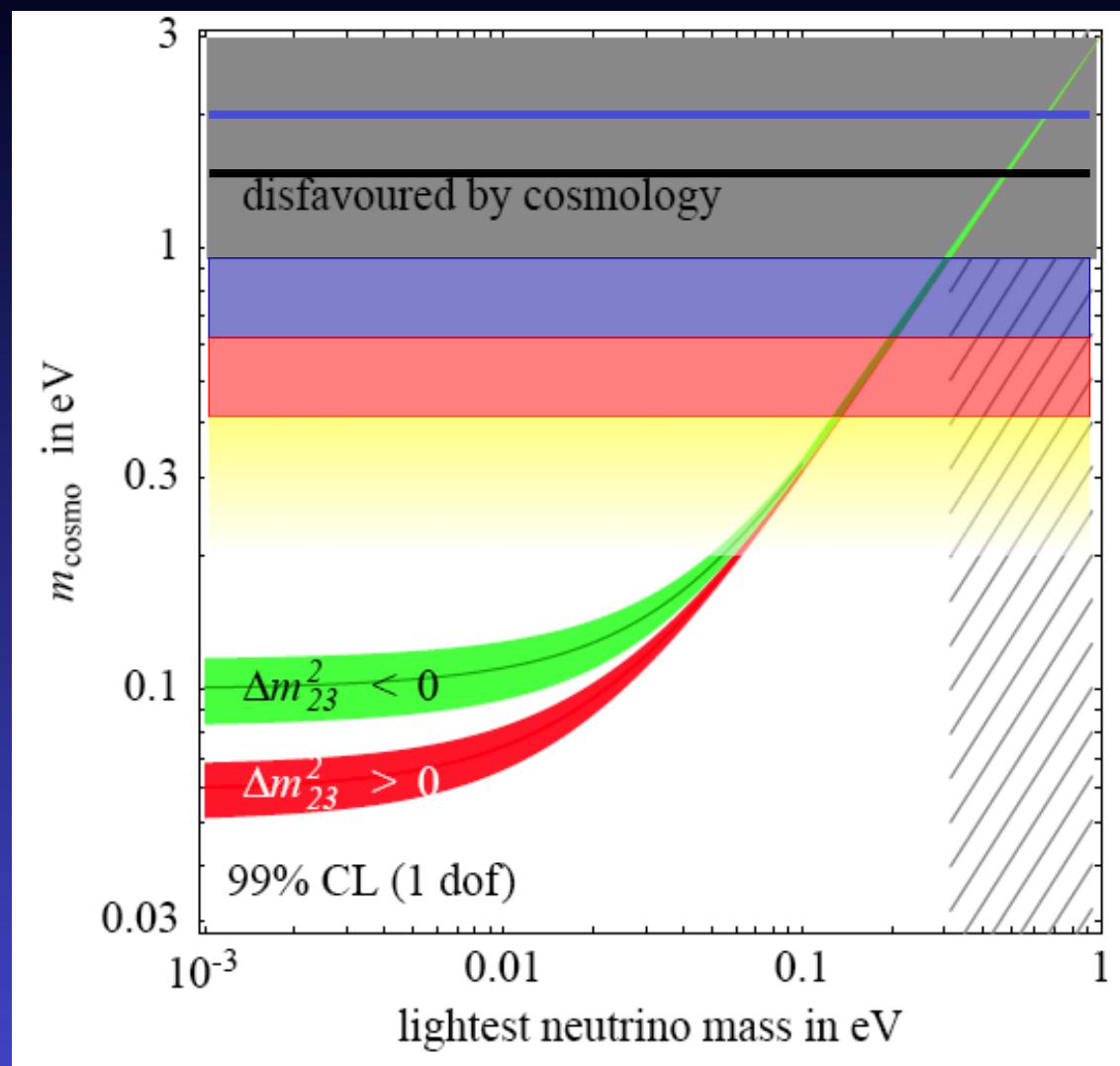
Cosmological data set	With WMAP3	$\Sigma$ bound ( $2\sigma$ )
WMAP		< 2.3 eV
WMAP + SDSS		< 1.2 eV
WMAP + SDSS + SN <sub>Riess</sub> + HST + BBN		< 0.78 eV
CMB + LSS + SN <sub>Astier</sub>		< 0.75 eV
CMB + LSS + SN <sub>Astier</sub> + BAO		< 0.58 eV
CMB + LSS + SN <sub>Astier</sub> + Ly- $\alpha$		< 0.21 eV
CMB + LSS + SN <sub>Astier</sub> + BAO + Ly- $\alpha$		< 0.17 eV

Fogli et al., PRD 75 (2007) 053001

Cosmological data set	With WMAP5	$\Sigma$ (at $2\sigma$ )
CMB		< 1.19 eV
CMB + HST + SN-Ia		< 0.75 eV
CMB + HST + SN-Ia + BAO		< 0.60 eV
CMB + HST + SN-Ia + BAO + Ly $\alpha$		< 0.19 eV

Fogli et al., PRD 78 (2008) 033010

# Neutrino masses in 3-neutrino schemes



- CMB
- CMB + galaxy clustering
- + HST, SNI-a...
- + BAO and/or bias
- + including Ly- $\alpha$

## Current cosmological bounds on neutrino masses

Dependence on the data set AND the cosmological model used.

Model	Observables	$\Sigma m_\nu$ (eV) 95% Bound
$\omega\text{CDM} + \Delta N_{\text{rel}} + m_\nu$	CMB+HO+SN+BAO	$\leq 1.5$
	CMB+HO+SN+LSSPS	$\leq 0.76$
$\Lambda\text{CDM} + m_\nu$	CMB+H0+SN+BAO	$\leq 0.61$
	CMB+H0+SN+LSSPS	$\leq 0.36$
$\Lambda\text{CDM} + m_\nu$	CMB (+SN)	$\leq 1.2$
$\Lambda\text{CDM} + m_\nu$	CMB+BAO	$\leq 0.75$
$\Lambda\text{CDM} + m_\nu$	CMB+LSSPS	$\leq 0.55$
$\Lambda\text{CDM} + m_\nu$	CMB+H0	$\leq 0.45$

González-García et al., JCAP 08 (2010) 117

# Absolute mass scale searches

Tritium  $\beta$   
decay

$$m_{\nu_e} = \left( \sum_i |U_{ei}|^2 m_i^2 \right)^{1/2}$$

2.2 eV

Neutrinoless  
double beta  
decay

$$m_{ee} = \left| \sum_i U_{ei}^2 m_i \right|$$

< 0.2-0.8 eV

Cosmology

$$\sim \sum_i m_i$$

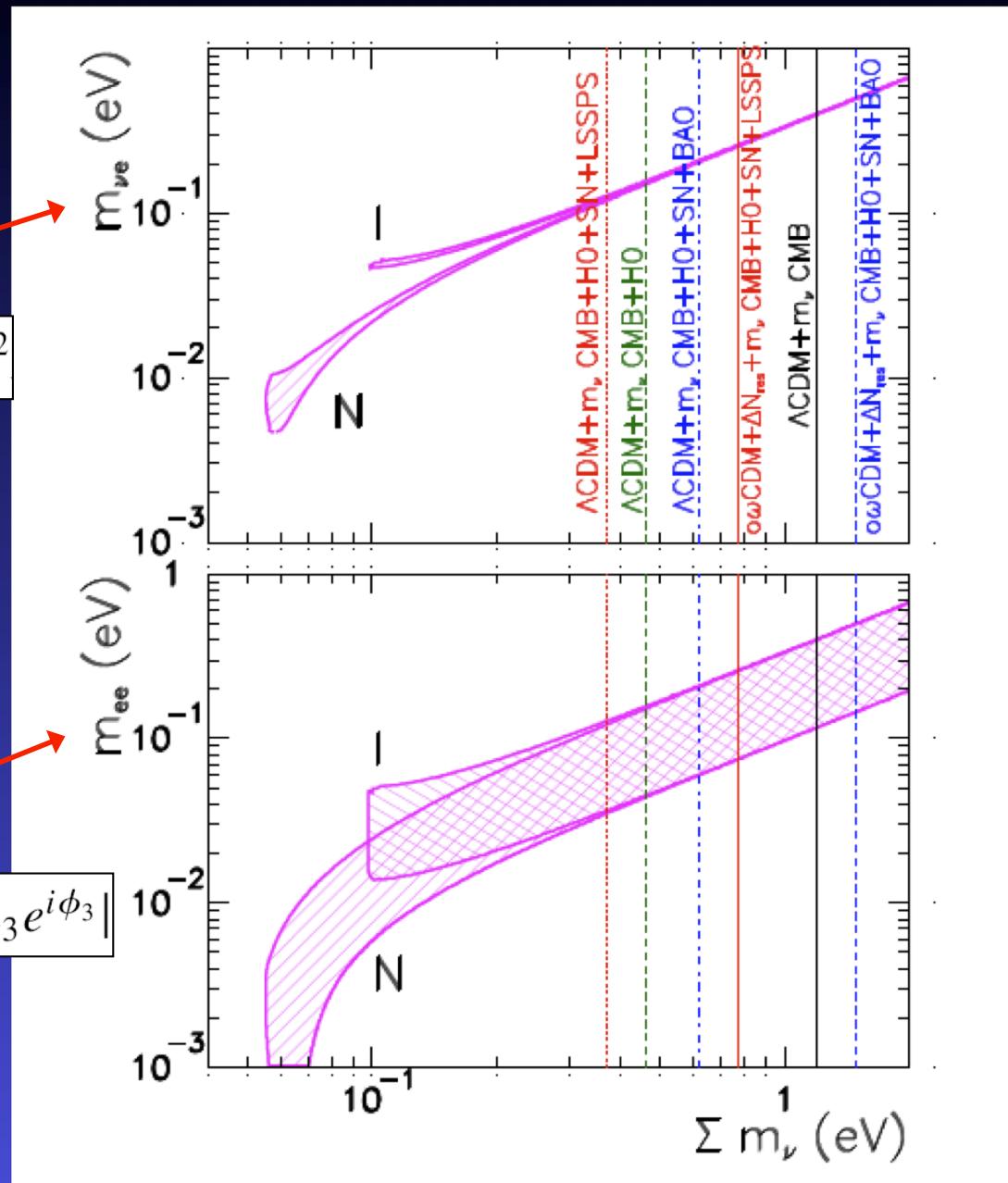
< 0.2-2.0 eV

# Tritium $\beta$ decay, $0\nu2\beta$ and Cosmology

$$[c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$

$$|c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

González-García et al.,  
JCAP 08 (2010) 117



# Non-standard relic neutrinos

The cosmological bounds on neutrino masses are modified if relic neutrinos have non-standard properties (or for non-standard models)

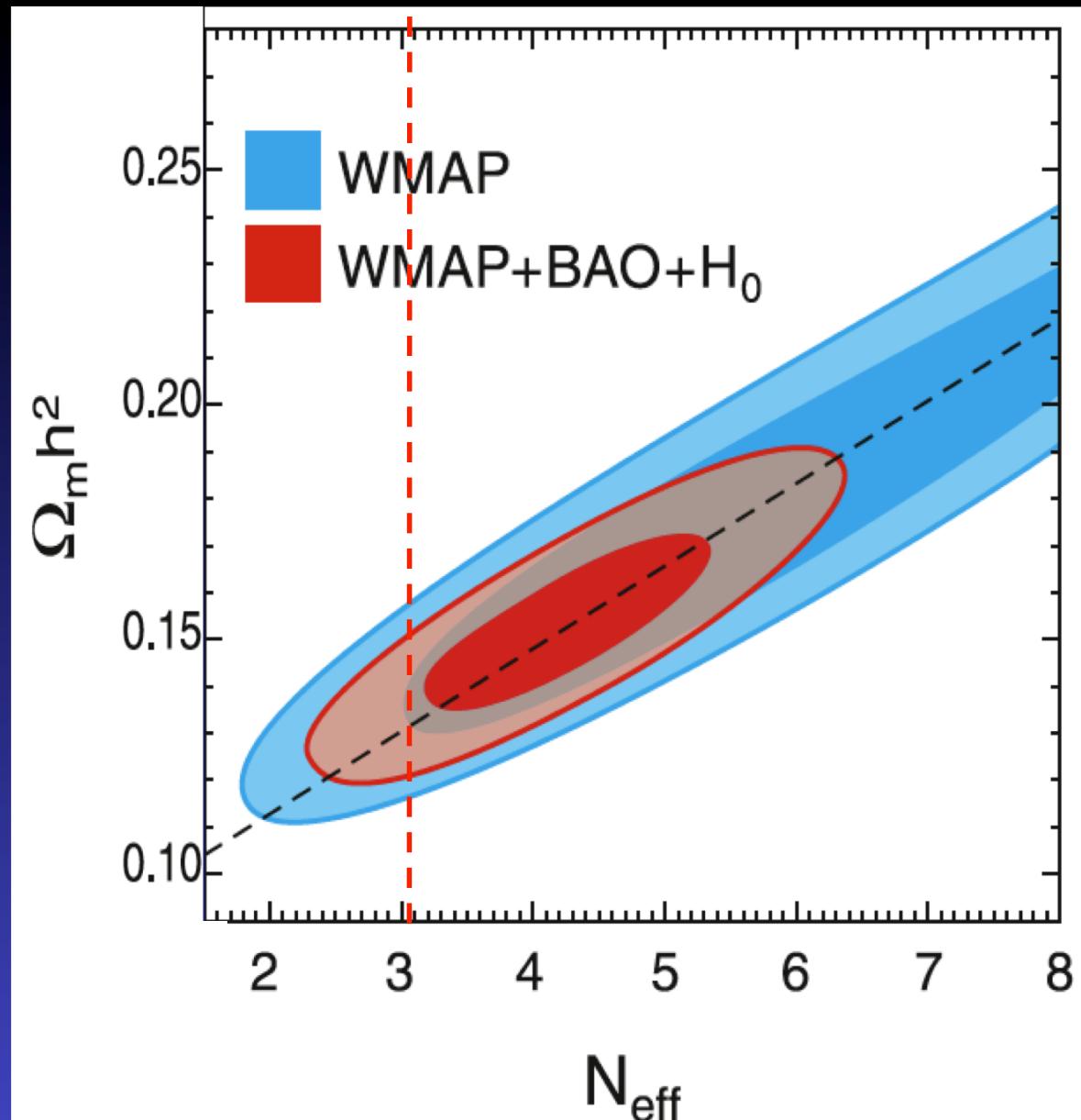
Two examples where the cosmological bounds do not apply

- Massive neutrinos strongly coupled to a light scalar field: they could annihilate when becoming NR
- Neutrinos coupled to the dark energy: the DE density is a function of the neutrino mass (mass-varying neutrinos)

Future sensitivities on  $m_\nu$  and  
 $N_{\text{eff}}$  from cosmology

# allowed range for $N_{\text{eff}}$

WMAP [7-year], arXiv:1001.4538



$2.7 < N_{\text{eff}} < 6.2$  (WMAP+BAO+ $H_0$ )

# Future bounds on $N_{\text{eff}}$

Forecast analysis:  
Bowen et al,  
MNRAS 334 (2002) 760



$$\sigma[N_{\text{eff}}] \sim 3 \text{ (WMAP)}$$

$$\sigma[N_{\text{eff}}] \sim 0.2 \text{ (Planck)}$$

Experiment	$f_{\text{sky}}$	$\theta_b$	$w_T^{-1/2}$	$w_p^{-1/2}$	$\Delta N_\nu$	$\Delta N_\nu$	$\Delta N_\nu$ (free $Y$ )
			[ $\mu \text{ K}'$ ]	[ $\mu \text{ K}'$ ]		TT	TT+TE+EE
Planck	0.8	7'	40	56	0.6	0.20	0.24
ACT	0.01	1.7'	3	4	1	0.47	0.9
ACT + Planck					0.4	0.18	0.24
CMBPOL	0.8	4'	1	1.4	0.12	0.05	0.09

Example of future  
CMB satellite

Bashinsky & Seljak, PRD 69 (2004) 083002

# Future bounds on $N_{\text{eff}}$

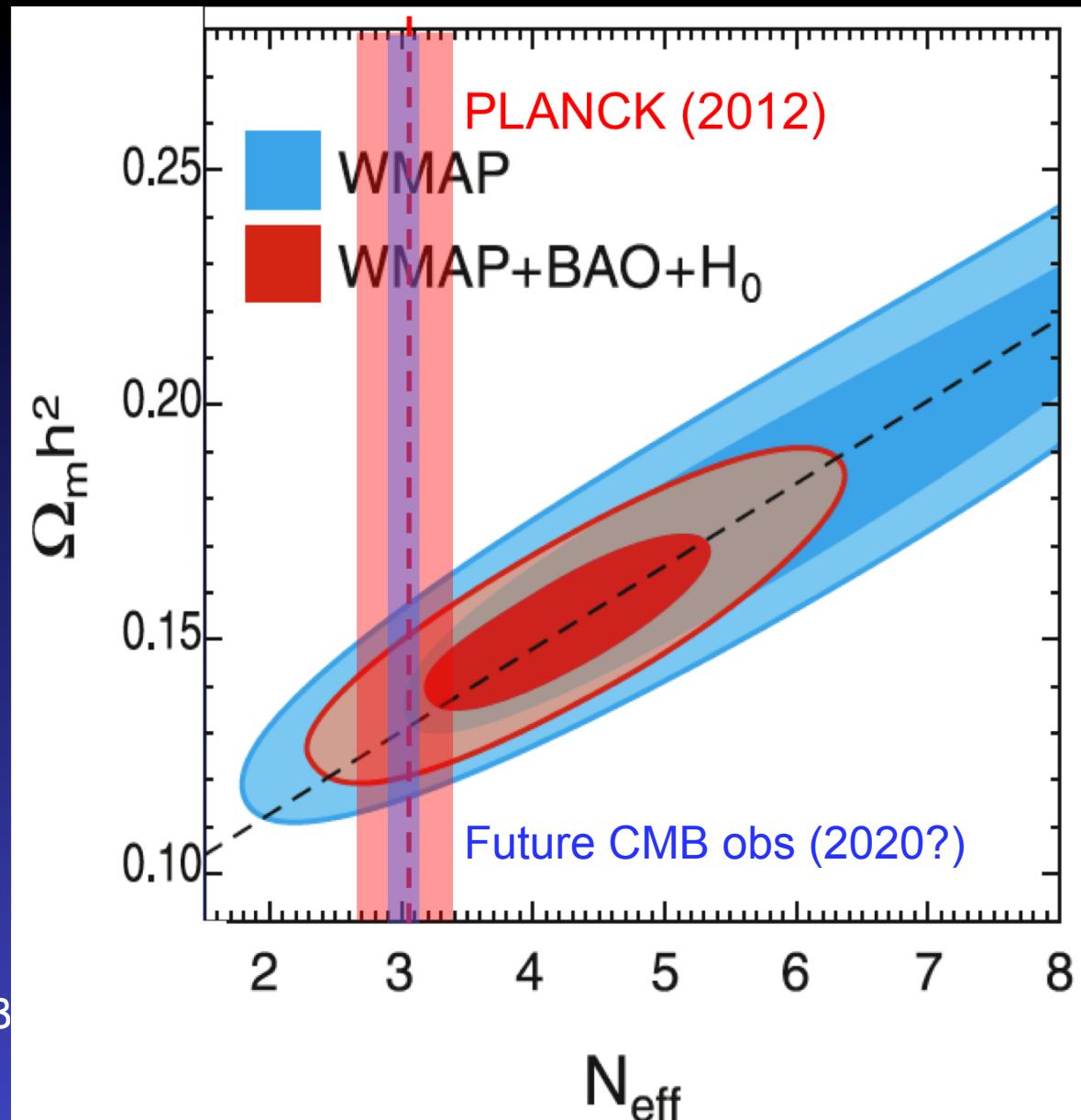
	Planck	P+BAO	P+HPS	P+HST	P+HST+BAO	P+HST+HPS
$\omega_{\text{dm}}$	0.22	0.24	0.20	0.21	0.21	0.19
$N_{\text{eff}}$	0.21	0.21	0.22	0.21	0.21	0.22
$\sum m_\nu$	0.68	0.81	0.44	0.67	0.73	0.44
$w$	2.14	1.16	0.72	0.74	0.76	0.55
$n_s$	0.46	0.48	0.49	0.46	0.48	0.48

**Table 3:** Projected sensitivity of Planck data (P) combined with LSS data to selected parameters of the  $\text{vanilla} + f_\nu + N_{\text{eff}} + w$  model. Given are the standard deviations of the marginalised posteriors, normalised to the values obtained with current CMB+HST+HPS data. Note that just like for current CMB data, the addition of BAO data shifts the posterior towards larger neutrino masses, resulting in a two-tailed pdf with a correspondingly larger standard deviation – this does not mean that the constraining power of Planck+BAO is worse than that of Planck alone. The marginalised posteriors of all the other parameters are very close to two-tailed Gaussians, and do not suffer from this effect.

Hamann et al, JCAP 07 (2010) 022

# allowed range for $N_{\text{eff}}$

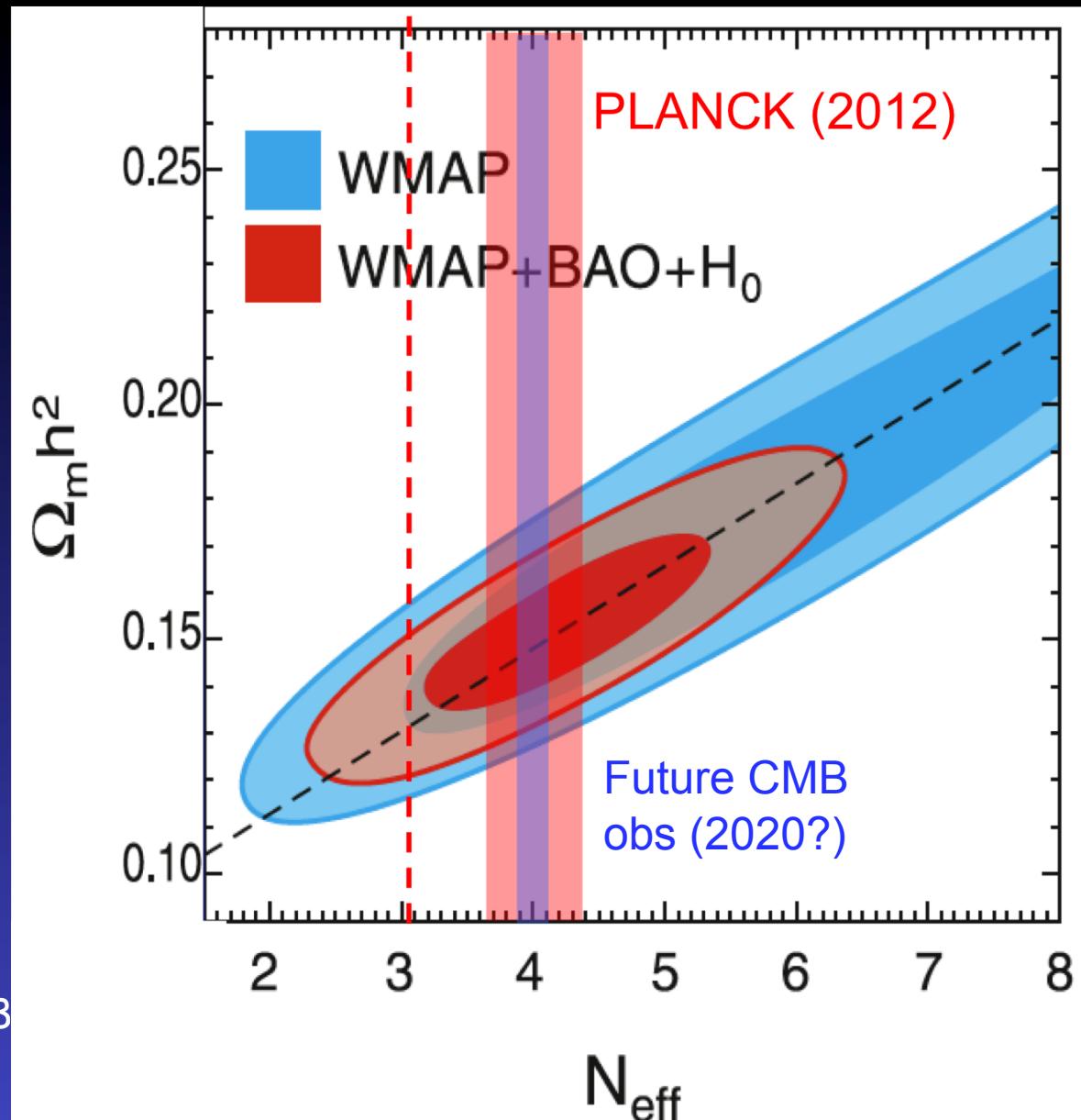
WMAP [7-year], arXiv:1001.4538



$2.7 < N_{\text{eff}} < 6.2$  (WMAP+BAO+H<sub>0</sub>)

# allowed range for $N_{\text{eff}}$

WMAP [7-year], arXiv:1001.4538



$2.7 < N_{\text{eff}} < 6.2$  (WMAP+BAO+ $H_0$ )

# Future sensitivities to $\Sigma m_\nu$

Future cosmological data will be available from

- o **CMB (Temperature & Polarization anis.)**
- o **High-z Galaxy redshift surveys**

Hannestad & Wong, JCAP 07 (2007) 004

Takada et al, PRD 73 (2006) 083520

- o **Galaxy cluster surveys**

Wang et al, PRL 95 (2005) 011302

- o **Weak lensing surveys (tomography)**

Hannestad et al, JCAP 06 (2006) 025

Song & Knox, PRD 70 (2004) 063510

- o **CMB lensing**

Perotto et al, JCAP 10 (2006) 013

Lesgourgues et al, PRD 73 (2006) 045021

- o **Fluctuations in the 21 cm H line**

Loeb & Wyithe, PRL 100 (2008) 161301

Pritchard & Pierpaoli, PRD 78 (2008) 065009

Forecasts indicate 0.01-0.15 eV sensitivities to  $\Sigma m_\nu$  are possible !!

# PLANCK+SDSS

- Fisher matrix analysis: expected sensitivities assuming a fiducial cosmological model, for future experiments with known specifications

Fiducial cosmological model:  
 $(\Omega_b h^2, \Omega_m h^2, h, n_s, \tau, \Sigma m_v) =$   
 $(0.0245, 0.148, 0.70, 0.98, 0.12, \Sigma m_v)$

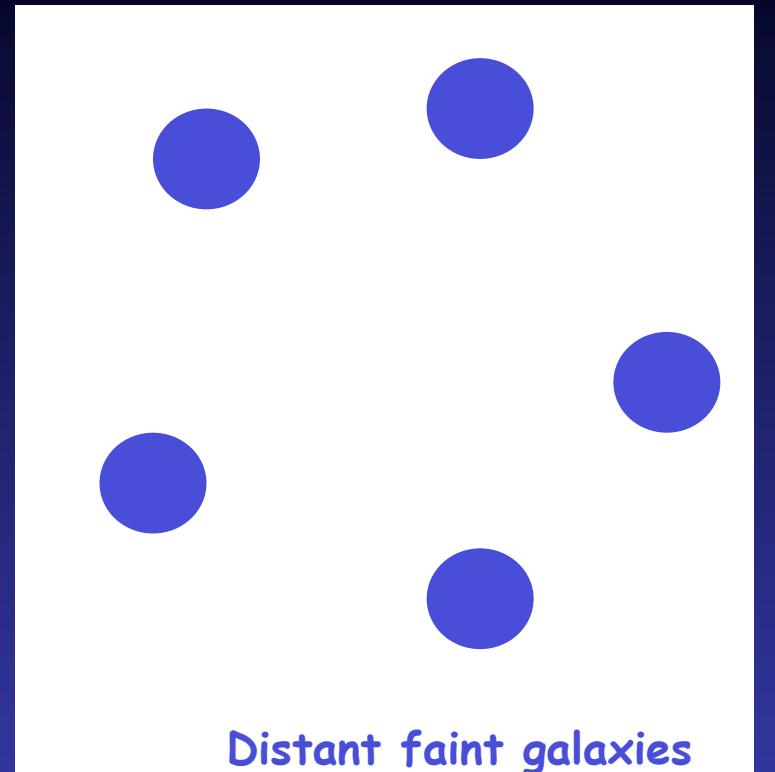
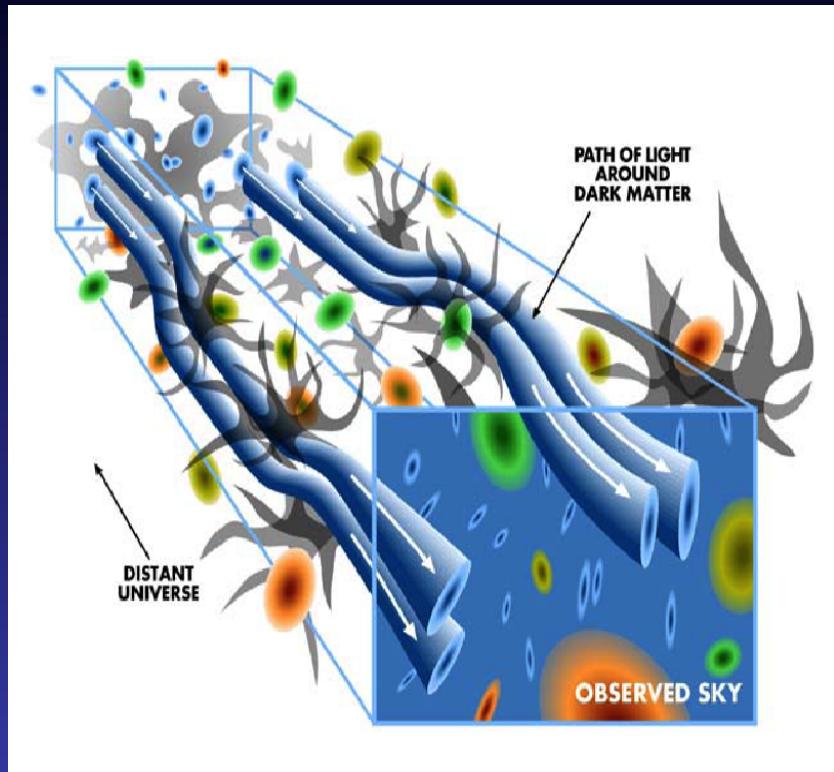
$\Sigma m$  detectable at  $2\sigma$  if  
larger than

0.21 eV (PLANCK+SDSS)

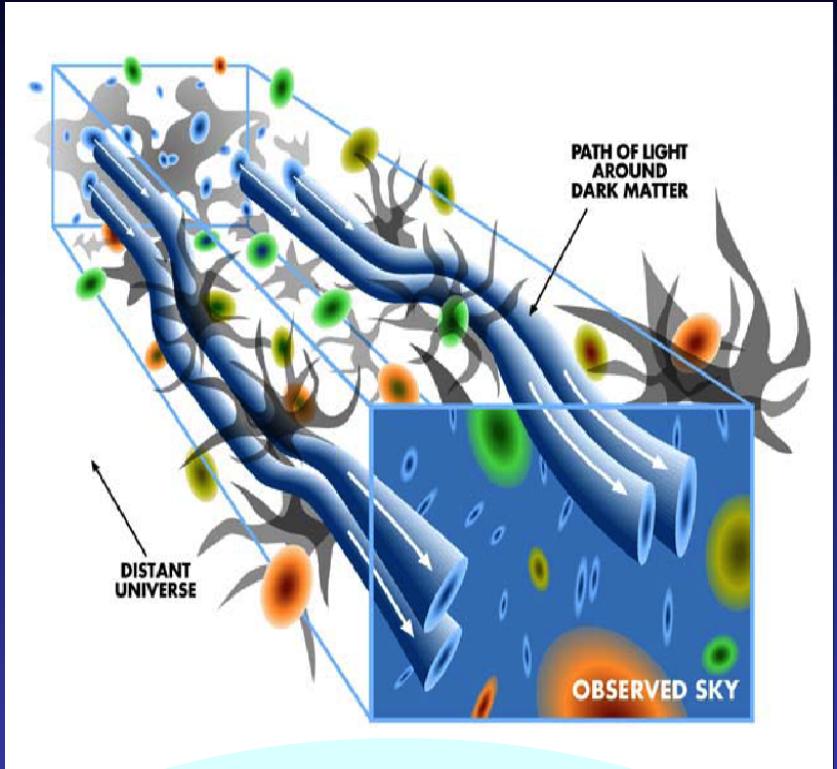
0.13 eV (CMBpol+SDSS)

Lesgourgues, SP & Perotto,  
PRD 70 (2004) 045016

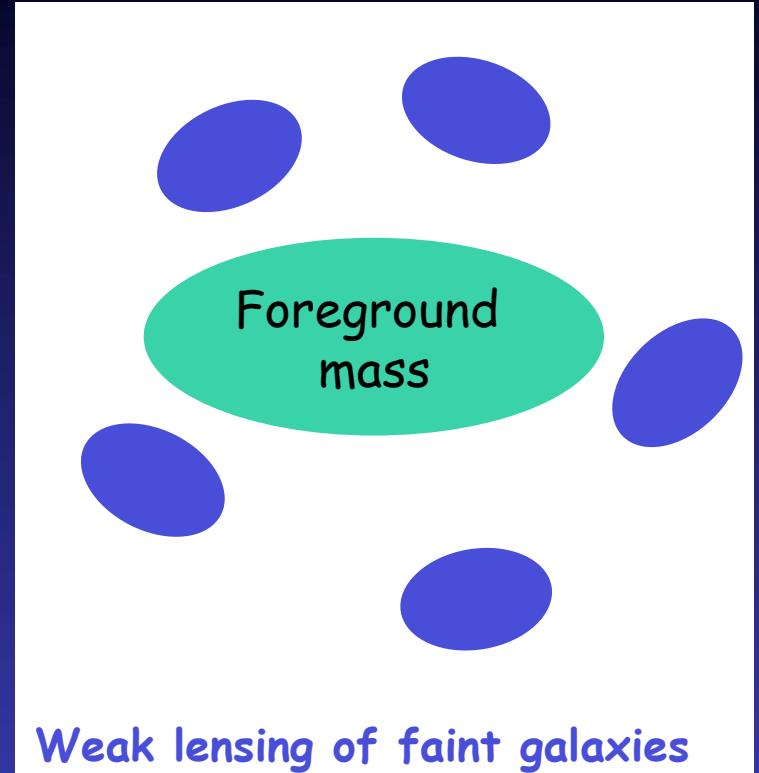
# Future sensitivities to $\Sigma m_\nu$ : weak gravitational lensing



# Future sensitivities to $\Sigma m_v$ : weak gravitational lensing



No bias uncertainty  
Small scales much closer  
to linear regime  
Tomography:  
3D reconstruction

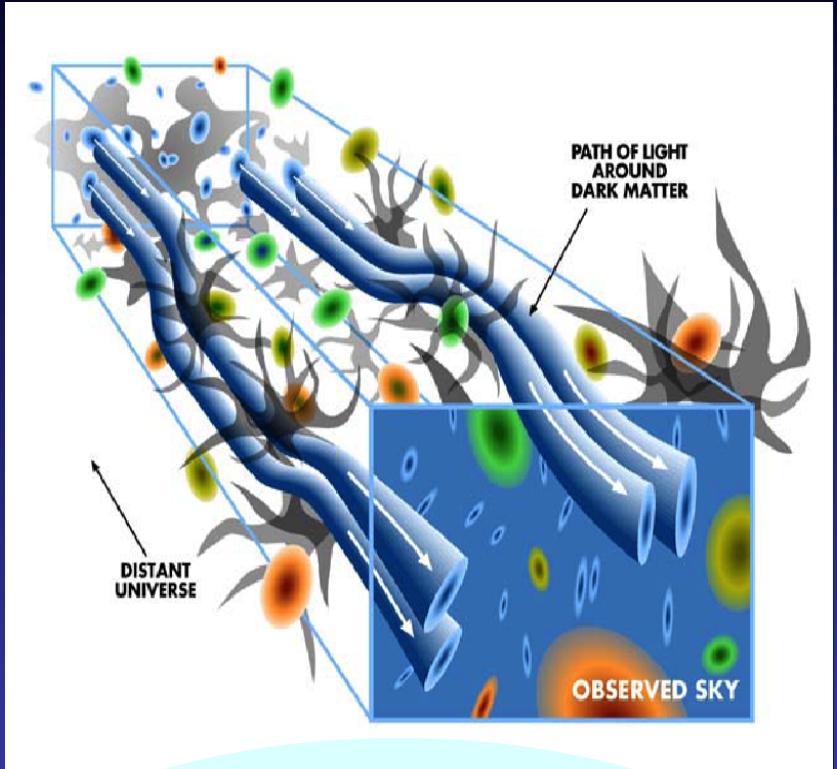


Weak lensing of faint galaxies

Frieman, Dodelson

Measure a large number  
of elliptically shaped galaxies

# Future sensitivities to $\Sigma m_\nu$ : weak gravitational lensing



No bias uncertainty  
Small scales much closer  
to linear regime  
Tomography:  
3D reconstruction

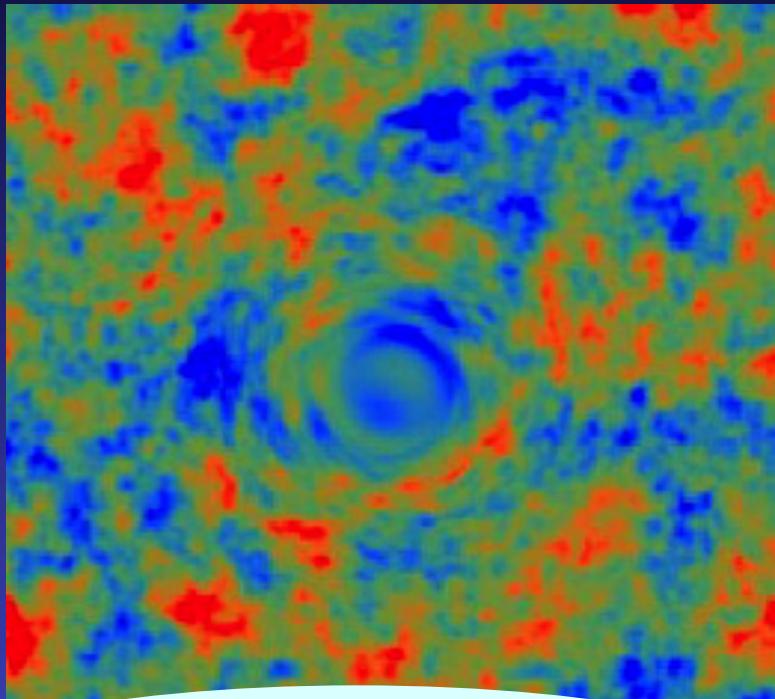
sensitivity of future  
weak lensing survey  
 $(4000^\circ)^2$  to  $m_\nu$

$$\sigma(m_\nu) \sim 0.1 \text{ eV}$$

Abazajian & Dodelson  
PRL 91 (2003) 041301

# Future sensitivities to $\Sigma m_\nu$ : weak gravitational lensing

lensing of the CMB signal



Makes CMB sensitive to  
smaller neutrino masses

sensitivity of CMB  
(primary + lensing)  
to  $m_\nu$

$$\sigma(m_\nu) = 0.15 \text{ eV (Planck)}$$

$$\sigma(m_\nu) = 0.044 \text{ eV (CMBpol)}$$

Kaplinghat, Knox & Song  
PRL 91 (2003) 241301

# CMB lensing: recent forecast analysis

$\sigma(M_v)$  in eV for future CMB experiments alone :

Lesgourges et al,  
PRD 73 (2006) 045021

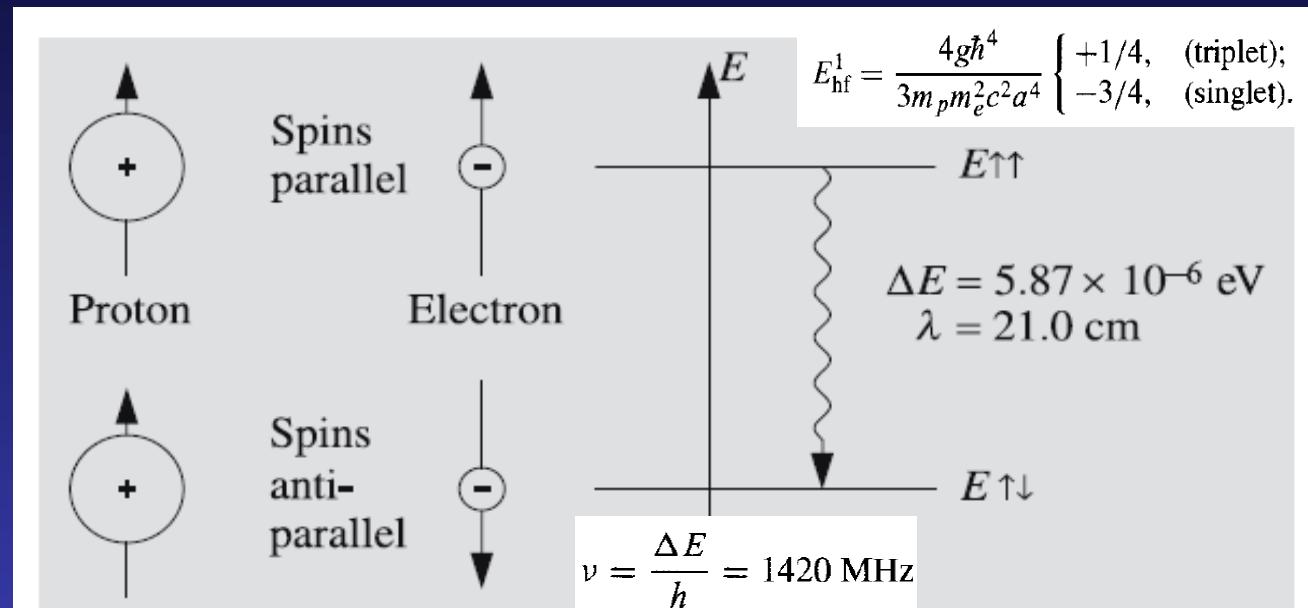
Free parameters:		8 parameters of minimal AMDM		
Lensing extraction:	no	no	yes	yes
Foreground cleaning:	perfect	none	perfect	none
QUaD+BICEP	1.3	1.6	0.31	0.36
BRAIN+CLOVER	1.5	1.8	0.34	0.43
PLANCK	0.45	0.49	0.13	0.14
SAMPAN	0.34	0.40	0.10	0.17
PLANCK+SAMPAN	0.32	0.36	0.08	0.10
Inflation Probe	0.14	0.16	0.032	0.036

Free parameters:		same + $\{\alpha, w, N_{\text{eff}}\}$		
Lensing extraction:	no	no	yes	yes
Foreground cleaning:	perfect	none	perfect	none
QUaD+BICEP	1.5	1.9	0.36	0.40
BRAIN+CLOVER	1.7	2.0	0.42	0.51
PLANCK	0.51	0.56	0.15	0.15
SAMPAN	0.37	0.44	0.12	0.18
PLANCK+SAMPAN	0.34	0.40	0.10	0.12
Inflation Probe	0.25	0.26	0.035	0.039

# “Measuring” even $m_\nu=0.05$ eV ?

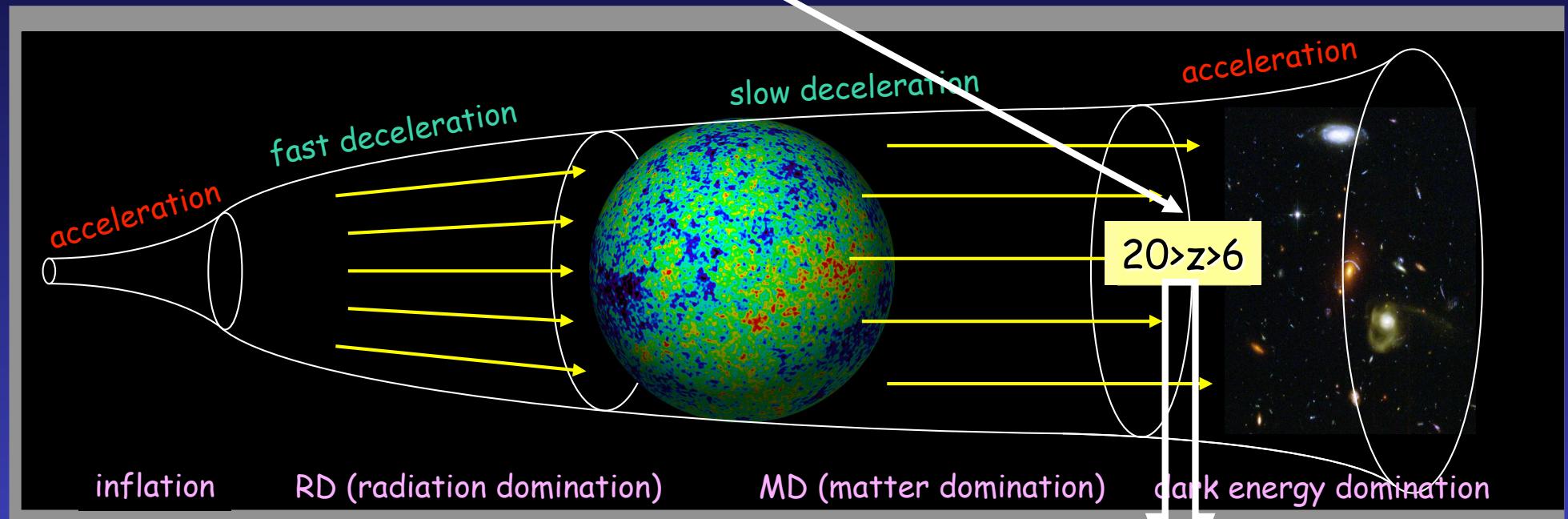
New cosmological observable as a potential probe of fluctuations at intermediate redshifts ( $6 < z < 20$ )  $\Rightarrow$  study of fluctuations in the 21cm line emitted by neutral H



**Fig. 5.8.** The origin of the hydrogen 21 cm line. The spins of the electron and the proton may be either parallel or opposite. The energy of the former state is slightly larger. The wavelength of a photon corresponding to a transition between these states is 21 cm

# “Measuring” even $m_\nu=0.05$ eV ?

New cosmological observable as a potential probe of fluctuations at intermediate redshifts ( $6 < z < 20$ )  $\Rightarrow$  study of fluctuations in the 21cm line emitted by neutral H



Redshifted line:  
2.1 m at  
redshift 10

$$\langle \tilde{\delta}_{21}(\mathbf{k}_1) \tilde{\delta}_{21}(\mathbf{k}_2) \rangle \equiv (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2) P_{21}(\mathbf{k}_1)$$

power spectrum of 21 cm brightness fluctuations  $P_{21}(k)$

# Future sensitivities on $\Sigma m_\nu$ from 21cm observ.

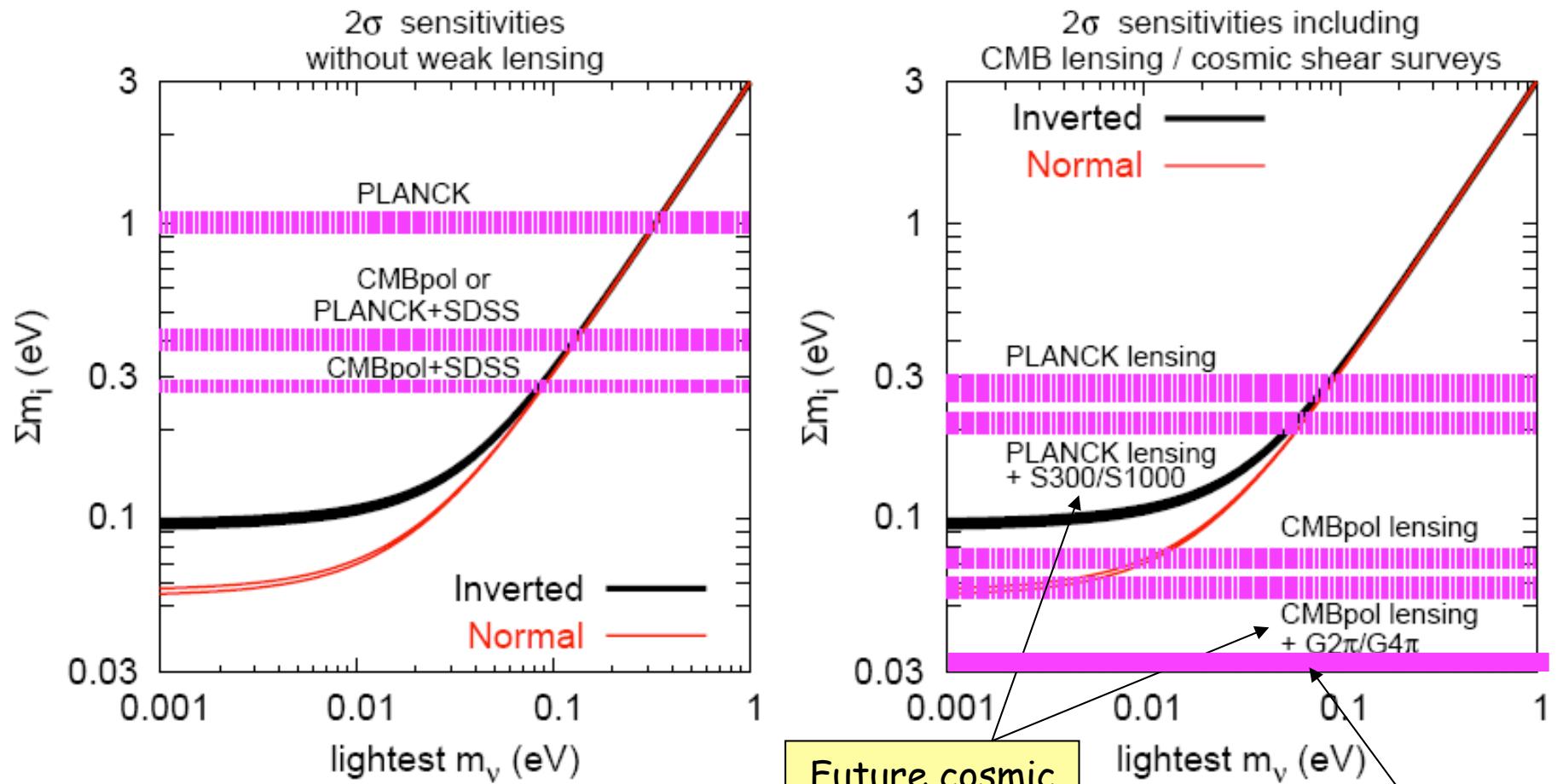
	$\Omega_m h^2$	$\Omega_b h^2$	$\Omega_\Lambda$	$w$	$n_s$	$A_s^2$	$\tau$	$Y_{He}$	$M_\nu$
Fiducial	0.147	0.023	0.7	-1	0.95	26.6	0.1	0.24	0.3
SDSS	0.456	0.083	0.117	1.21	0.503	$\infty$	-	-	6.16
G1	0.119	0.0207	0.0358	0.574	0.174	$\infty$	-	-	1.28
G2	0.0354	0.00593	0.295	1.22	0.0482	$\infty$	-	-	1.01
G3	0.0252	0.00438	0.0076	0.122	0.037	$\infty$	-	-	0.272
MWA	0.0317	0.00761	0.972	3.13	0.0487	$\infty$	-	-	0.749
SKA	0.00191	0.00056	0.234	0.747	0.0054	$\infty$	-	-	0.175
FFTT	4.94e-05	4.77e-05	0.0045	0.014	0.0002	$\infty$	-	-	0.009
Planck	0.0045	0.00024	0.068	0.18	0.0074	0.26	0.0048	0.011	0.38
+SDSS	0.0033	0.00024	0.023	0.11	0.0074	0.254	0.0046	0.0103	0.272
+G1	0.0016	0.00021	0.013	0.081	0.0068	0.245	0.0044	0.01	0.136
+G2	0.00089	0.00022	0.037	0.149	0.0067	0.243	0.0044	0.0099	0.104
+G3	0.00051	0.00016	0.003	0.021	0.0051	0.24	0.0043	0.0081	0.052
+MWA	0.00146	0.00021	0.053	0.17	0.0066	0.242	0.0044	0.0101	0.144
+SKA	0.00029	0.00014	0.020	0.065	0.003	0.236	0.0043	0.0044	0.080
+FFTT	4.23e-05	3.97e-05	0.004	0.011	0.0002	0.23	0.0043	0.0030	0.0075
CosmicVar	0.00244	4.16e-05	0.030	0.033	0.0024	0.124	0.0023	0.0028	0.222
+FFTT	3.3e-05	2.58e-05	0.0024	0.0076	0.0002	0.111	0.0021	0.0011	0.0068

Future  
Low- $\nu$  radio  
telescopes

Pritchard & Pierpaoli, PRD 78 (2008) 065009

[also Loeb & Wyithe, PRL 100 (2008) 161301; Mao et al, PRD 78 (2008) 023529]

# Summary of future sensitivities



# Summary of future sensitivities

Probe	Potential sensitivity (short term)	Potential sensitivity (long term)
CMB	0.4-0.6	0.4
CMB with lensing	0.1-0.15	0.04
CMB + Galaxy Distribution	0.2	0.05-0.1
CMB + Lensing of Galaxies	0.1	0.03-0.04
CMB + Lyman- $\alpha$	0.1-0.2	Unknown
CMB + Galaxy Clusters	-	0.05
CMB + 21 cm	-	0.0003-0.1

**Table 1.** Future probes of neutrino mass, as well as their projected sensitivity to neutrino mass. Sensitivity in the short term means achievable in approximately 5-7 years, while long term means 7-15 years.

Hannestad, arXiv:1007.0658

**End of 2nd lecture**



# Conclusions

Cosmological observables can be used to bound (or measure) neutrino properties, in particular the sum of neutrino masses (info complementary to laboratory results)

The radiation content of the Universe ( $N_{\text{eff}}$ ) will be very constrained in the near future (Planck)

Current bounds on the sum of neutrino masses from cosmological data  
(best  $\sum m_\nu < 0.4\text{-}0.6 \text{ eV}$ , conservative  $\sum m_\nu < 1 \text{ eV}$ )

Different cosmological observations in the next future → Sub-eV sensitivity (0.1-0.2 eV and better) → Test degenerate mass region and eventually the minimum total mass for IH case